On the second Painlevé equation and its higher analogues

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- Painlevé equations
- Integrability
- Symmetries
- 4 Sigma-coordinates
- 6 Higher analogues

Example 1.



Definition 1. If singular points of the solution of the differential equation depend on the initial data, then such points are called *movable*.

The Painlevé property

Definition 2. A differential equation satisfies *the Painlevé property*, if its explicit solution has movable simple poles only.

Problem 1. Classify all equations satisfied the Painlevé property in the given form

$$\frac{d^2w}{dz^2} = P\left(z, w, \frac{dw}{dz}\right),\tag{1}$$

where P(z, w, w') is a meromorphic function in z and a rational function in w and w'.

Example 1.

No singularity pointsSimple poleBranch singularity points
$$\frac{dw}{dz} = w$$
 $\frac{dw}{dz} = w^2$ $\frac{dw}{dz} = \frac{1}{2}w^3$ $w(z) = z_0 e^z$ $w(z) = \frac{1}{z_0 - z}$ $w(z) = \pm \frac{1}{\sqrt{z_0 - z}}$

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The Painlevé property and isomonodromic problem

$$\begin{cases} \left(\frac{d}{d\lambda} + A(\lambda, z)\right) Y(\lambda, z) = 0, & \det Y(\lambda, z) = 1; \\ \left(\frac{d}{dz} + B(\lambda, z)\right) Y(\lambda, z) = 0, & \operatorname{tr} B(\lambda, z) = 0, B(\lambda, z) = -\frac{d}{dt} Y \cdot Y^{-1}. \end{cases}$$

the compatibility condition $\Longrightarrow \frac{d}{dz}A = \frac{d}{d\lambda}B + [B, A]$, tr B = 0. (2)



The confluence scheme for the Painlevé equations

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The confluence scheme for the Painlevé equations

PVI
$$[w(z); \alpha, \beta, \gamma, \delta]; (0, 0, 0, 0)_{(0,1,\infty,z)}$$

$$w'' = \frac{1}{2} \left(\frac{1}{w} + \frac{1}{w-1} + \frac{1}{w-z} \right) (w')^2 - \left(\frac{1}{z} + \frac{1}{z-1} + \frac{1}{w-z} \right) w' + \frac{w(w-1)(w-z)}{z^2(z-1)^2} \left(\alpha + \beta \frac{z}{w^2} + \gamma \frac{z-1}{(w-1)^2} + \delta \frac{z(z-1)}{(w-z)^2} \right).$$

.

PII
$$[w(z); \alpha]; (3)_{(\infty)}$$

 $w'' = 2w^3 + wz + \alpha.$
PI $[w(z)]; (5/2)_{(\infty)}$
 $w'' = 6w^2 + z.$

Example 2. (from P_{II} to P_{I})

 $w'' = 2w^3 + wz + \alpha \quad \left\{ w \mapsto \varepsilon w + \varepsilon^{-5}, \, z \mapsto \varepsilon^2 z - 6\varepsilon^{-10}, \, \alpha \mapsto 4\varepsilon^{-15} \right\}$ $w'' = 2\varepsilon^6 w^3 + 6w^2 + \varepsilon^6 z w + z \quad \left\{ \varepsilon \to 0 \right\} \quad w'' = 6w^2 + z$

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The Hamiltonian structure

$$\mathsf{P}_{\mathsf{II}}[w(z);b]: \quad w'' = 2w^3 + zw + \left(b - \frac{1}{2}\right). \tag{3}$$

Okamoto variables

$$q = w, \quad p = w' + w^2 + \frac{z}{2};$$
 $\{q, p\} = 1, \quad \{q, q\} = \{p, p\} = 0.$
(4)

$$H_{II}[b] = \frac{1}{2}p^{2} - \left(q^{2} + \frac{z}{2}\right)p - bq \quad \Leftrightarrow \quad \begin{cases} q' = \frac{\partial H}{\partial p} = p - q^{2} - \frac{z}{2}, \\ p' = -\frac{\partial H}{\partial q} = 2pq + b. \end{cases}$$
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$$\begin{split} A(\lambda, z) &= \lambda^2 \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \lambda \begin{pmatrix} 0 & u \\ -2u^{-1}v & 0 \end{pmatrix} + \begin{pmatrix} v + z/2 & -uw \\ -2u^{-1}(wv - b) & -v - z/2 \end{pmatrix}, \\ B(\lambda, z) &= \frac{\lambda}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 0 & u \\ -2u^{-1}v & 0 \end{pmatrix}. \\ \left\{ \frac{d \ln u}{dz} = -w, \frac{dw}{dz} = w^2 + v + \frac{z}{2}, \frac{dv}{dz} = -2wv + b \right\} \implies \mathsf{P}_{\mathsf{H}} \left[w(z); b - \frac{1}{2} \right]. \end{split}$$

Remark.
$$P_{34}[v(z); b]: \frac{d^2v}{dz^2} = \frac{1}{2z} \left(\frac{dv}{dz}\right)^2 - 2v^2 - zv - \frac{b^2}{2v}$$

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Definition 3.

- A Bäcklund transformation transforms one PDE solution to another its solution.
- An auto-Bäcklund transformation leaves a PDE invariant.

Example 3. $w \mapsto -w$: $-w'' = -2w^3 - zw + b \Rightarrow P_{II}[w; b] \mapsto P_{II}[w; -b]$.

Claim 1.

The transformation s in the given form preserves the solution of the P_{II} equation

$$s\left(\tilde{q},\tilde{p},\tilde{z};\tilde{b}\right) = s\left(q + \frac{b}{p}, p, z; -b\right),\tag{6}$$

The coordinates q and p are canonical.

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The transformation r in the given form preserves the solution of the P_{II} equation

$$r\left(\hat{q},\hat{p},\hat{z};\hat{b}\right) = r\left(-q,-p+2q^{2}+z,z;1-b\right).$$
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Theorem 1. The Bäcklund transformation for the second Painlevé equation is the following one

$$\tilde{w} = w - \frac{1}{2} \frac{2\alpha - \varepsilon}{\varepsilon w' - w^2 - \frac{z}{2}}, \quad \alpha = 1 - \tilde{\alpha}, \quad \varepsilon = \pm 1.$$
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Remark.

 $\varepsilon = 1$, then (8) coincides with the transformation *s*; $\varepsilon = -1$, then (8) coincides with the composition $r \circ s$.

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$$T_k = (rs)^k, \quad S_k = (rs)^k s, \quad k \in \mathbb{Z}.$$
 (9)

k	 -2	-1	0	1	2	
T_k	 srsr	sr	1	rs	rsrs	
S_k	 srsrs	srs	5	r	rsr	

Compositions of Bäcklund transformations



Claim 3.

1 $T_k(b) = b - k, S_k(b) = k - b, k \in \mathbb{Z}$. **2** $s^2 = 1$ and $r^2 = 1$.

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The affine Weyl group and its extension

Theorem 2. Bäcklund transformations r and s are generators of the affine Weyl group of type $A_1^{(1)}$,

$$W = \{T_k, S_k : k \in \mathbb{Z}\} = \langle r, s \rangle, \tag{10}$$

with fundamental relations $r^2 = 1$, $s^2 = 1$.

Remark. If we set $s_0 = s$, $s_1 = rsr$, and $\pi = r$, we obtain the extension of W:

$$\tilde{W} = \langle s_0, s_1; \pi \rangle, \quad s_0^2 = 1, \quad s_1^2 = 1, \quad \pi s_0 = s_1 \pi.$$
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The Cartan matrix The Dynkin diagram

The discrete dynamics d-P $\left(A_1^{(1)}/E_7^{(1)}\right)$ [KNY17]:

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The discrete dynamics d-P $\left(A_1^{(1)}/E_7^{(1)}\right)$ [KNY17]:

$$\begin{cases} \overline{p} = -p + 2q^2 + z, \\ \overline{q} = -q + \frac{1-b}{\overline{p}}. \\ \end{cases}$$
(12)

Generating rational solutions

The particular rational solution

$$q, p; b) = \left(0, \frac{z}{2}; \frac{1}{2}\right) \tag{13}$$

	••••	T_{-1}	T_0	T_1	 	S_{-1}	S_0	S_1	
q		$-\frac{1}{z}$	0	$\frac{1}{z}$	 	$\frac{2z^3-4}{z\left(z^3+4\right)}$	$\frac{1}{z}$	0	
p		$\frac{z^3+4}{2z^2}$	$\frac{z}{2}$	$\frac{z}{2}$	 	$\frac{z^3+4}{2z^2}$	$\frac{z}{2}$	$\frac{z}{2}$	
b		$\frac{3}{2}$	$\frac{1}{2}$	$-\frac{1}{2}$	 	$-\frac{3}{2}$	$-\frac{1}{2}$	$\frac{1}{2}$	

Claim 4. If m = 2b - k holds, rational solutions obtained by actions of T_k and S_m on $(q, p; b) = \left(0, \frac{z}{2}; \frac{1}{2}\right)$ are the same.

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The second Painlevé equation

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p	 $\frac{z^3+4}{2z^2}$	$\frac{z}{2}$	$\frac{z}{2}$		 $\frac{z^3+4}{2z^2}$	$\frac{z}{2}$	$\frac{z}{2}$	
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The second Painlevé equation

Sigma-coordinates

$$\sigma\left(z\right) = \mathsf{H}_{\mathsf{H}}\left(b\right),\tag{14}$$

$$\sigma' = -\frac{1}{2}p, \quad \sigma'' = -pq - \frac{1}{2}b.$$
 (15)

$$(\sigma'')^{2} + 4(\sigma')^{3} + 2\sigma'(z\sigma' - \sigma) - \frac{1}{4}b^{2} = 0.$$
(16)

Claim 5. Sigma – coordinates are
$$\log -$$
 symplectic
 $\Omega = d (\ln \sigma') \wedge d\sigma''.$

The mKdV equation: some symmetries

$$v_t - 6v^2 v_x + v_{xxx} = 0. (17)$$

Remark. The mKdV equation is the result of the Bäcklund transformation of the KdV equation by the Miura transformation.

$$\begin{array}{c|c} u_t + 6uu_x + u_{xxx} = 0 \Rightarrow \\ \text{the KdV equation} \end{array} \Rightarrow \begin{array}{c|c} u = v_x - v^2 \\ \text{the Miura transformation} \end{array} \Rightarrow \begin{array}{c} v_t - 6v^2v_x + v_{xxx} = 0. \\ \text{the mKdV equation} \end{array}$$

The symmetry

$$X = x\partial_x + 3t\partial_t - v\partial_v$$
Independent coordinates

$$v(x,t) = \frac{w(z(x,t))}{(3t)^{1/3}},$$

$$z(x,t) = \frac{x}{(3t)^{1/3}}$$

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The mKdV equation: some symmetries

$$v_t - 6v^2 v_x + v_{xxx} = 0. (17)$$

Remark. The mKdV equation is the result of the Bäcklund transformation of the KdV equation by the Miura transformation.

$$\begin{array}{c|c} u_t + 6uu_x + u_{xxx} = 0 \Rightarrow \\ \text{the KdV equation} \end{array} \Rightarrow \begin{array}{c|c} u = v_x - v^2 \\ \text{the Miura transformation} \end{array} \Rightarrow \begin{array}{c|c} v_t - 6v^2v_x + v_{xxx} = 0. \\ \text{the mKdV equation} \end{array}$$

The symmetry

$$X = x\partial_x + 3t\partial_t - v\partial_v$$
Independent coordinates

$$v(x, t) = \frac{w(z(x, t))}{(3t)^{1/3}},$$

$$z(x, t) = \frac{x}{(3t)^{1/3}}$$

$$\mathsf{P}_{\mathsf{II}}\left[w\left(z
ight);lpha_{1}
ight]:w^{\prime\prime}=2w^{3}+zw+lpha_{1},\quad lpha_{1}=\mathsf{const}\,.$$

The stationary mKdV hierarchy

The stationary KdV hierarchy

$$\partial_{t_{2n+1}}u + \partial_x \ell_{n+1} \left[u \right] = 0, \quad n \ge 0, \tag{18}$$

$$\partial_{x}\ell_{n+1} = \left(\partial_{x}^{3} + 4u\partial_{x} + 2u_{x}\right)\ell_{n}, \quad \ell_{0} = \frac{1}{2},$$
(19)

where ℓ_n is the Lenard operator.

Example 4. (n = 2) $u_{t_5} + u_x^{(5)} + 10uu_{xxx} + 20u_xu_{xx} + 30u^2u_x = 0$ \downarrow the Miura transformation

The stationary mKdV hierarchy [Jos04] $\partial_{t_{2n+1}}v + (\partial_{xx} + 2\partial_{x}v) \ell_{n} [v_{x} - v^{2}] = 0, \quad n \ge 0, \quad (20)$ $\partial_{x}\ell_{n+1} = (\partial_{x}^{3} + 4 (v_{x} - v^{2}) \partial_{x} + 2 (v_{x} - v^{2})_{x}) \ell_{n}, \quad \ell_{0} = \frac{1}{2}. \quad (21)$

Example 5. (n = 2) $v_{t_5} + v_x^{(5)} - 10v^2v_{xxx} - 40v_xv_{xx} - 10v_x^3 + 30v^4x_x = 0$

The stationary mKdV hierarchy

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The stationary PII hierarchy

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The symmetry

$$X_n = x\partial_x + (2n+1)\partial_{t_{2n+1}} - v\partial_v$$
Independent coordinates

$$v(x, t_{2n+1}) = \frac{w(z(x, t))}{((2n+1)t_{2n+1})^{\frac{1}{2n+1}}},$$

$$z(x, t) = \frac{x}{((2n+1)t_{2n+1})^{\frac{1}{2n+1}}},$$

The stationary PII hierarchy [Jos04]

$$\left(\frac{d}{dz} + 2w\right)\mathcal{L}_n\left[w' - w^2\right] = zw + \alpha_n, \quad n \ge 1,$$
(23)

$$\partial_{z}\mathcal{L}_{n+1} = \left(\partial_{z}^{3} + 4\left(w' - w^{2}\right)\partial_{z} + 2\left(w' - w^{2}\right)'\right)\mathcal{L}_{n}, \quad \mathcal{L}_{1} = w' - w^{2}.$$
(24)

Example 6. (n = 0) $(z - 1) w = -\alpha_0$ **Example 7.** (n = 2) $w''' - 10w^2w'' - 10ww'^2 + 6w^5 = zw + \alpha_2$

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(22)

The stationary PII hierarchy

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Independent coordinates

$$v(x, t_{2n+1}) = \frac{w(z(x, t))}{((2n+1)t_{2n+1})^{\frac{1}{2n+1}}},$$

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$$\partial_{t_{2n+1}} v + (\partial_{xx} + 2\partial_x v) \ell_n \left[v_x - v^2 \right] = 0, \quad n \ge 0.$$

$$(22)$$

The symmetry

$$X_{n} = x\partial_{x} + (2n+1)\partial_{t_{2n+1}} - v\partial_{v}$$
Independent coordinates

$$v(x, t_{2n+1}) = \frac{w(z(x, t))}{((2n+1)t_{2n+1})^{\frac{1}{2n+1}}},$$

$$z(x, t) = \frac{x}{((2n+1)t_{2n+1})^{\frac{1}{2n+1}}},$$

The stationary PII hierarchy [Jos04]

$$\left(\frac{d}{dz}+2w\right)\mathcal{L}_n\left[w'-w^2\right]=zw+\alpha_n,\quad n\ge 1,$$
(23)

$$\partial_{z}\mathcal{L}_{n+1} = \left(\partial_{z}^{3} + 4\left(w' - w^{2}\right)\partial_{z} + 2\left(w' - w^{2}\right)'\right)\mathcal{L}_{n}, \quad \mathcal{L}_{1} = w' - w^{2}.$$
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The non-stationary PII hierarchy

The Virasoro symmetry $T_{n} = \sum_{k=0}^{n} (2k+1) T_{k+1} \partial_{T_{k+1}} \begin{cases} t_{0} = -z, \\ t_{k} = \frac{(2k+1) T_{k+1}}{((2n+1) T_{n+1})^{\frac{2k+1}{2n+1}}} \end{cases}$

The non-stationary PII hierarchy [MM07] $\left(\frac{d}{dz} + 2w\right) \left(\mathcal{L}_n \left[w' - w^2\right] + \sum_{k=1}^{n-1} t_k \mathcal{L}_k \left[w' - w^2\right]\right) = zw + \alpha_n, \quad n \ge 1, \quad (25)$ $\partial_z \mathcal{L}_{n+1} = \left(\partial_z^3 + 4\left(w' - w^2\right)\partial_z + 2\left(w' - w^2\right)'\right)\mathcal{L}_n, \quad \mathcal{L}_1 = w' - w^2. \quad (26)$

Example 8.

 $(n=2) \qquad \left(w'''' - 10w^2w'' - 10ww'^2 + 6w^5\right) + t_1\left(w'' - 2w^3\right) = zw + \alpha_2$

The non-stationary PII hierarchy

The Virasoro symmetry Independent coordinates

$$T_{n} = \sum_{k=0}^{n} (2k+1) T_{k+1} \partial_{T_{k+1}} \begin{vmatrix} t_{0} = -z, \\ t_{k} = \frac{(2k+1) T_{k+1}}{((2n+1) T_{n+1})^{\frac{2k+1}{2n+1}}} \end{vmatrix}$$

The non-stationary PII hierarchy [MM07] $\left(\frac{d}{dz} + 2w\right) \left(\mathcal{L}_n \left[w' - w^2\right] + \sum_{k=1}^{n-1} t_k \mathcal{L}_k \left[w' - w^2\right]\right) = zw + \alpha_n, \quad n \ge 1, \quad (25)$ $\partial_z \mathcal{L}_{n+1} = \left(\partial_z^3 + 4 \left(w' - w^2\right) \partial_z + 2 \left(w' - w^2\right)'\right) \mathcal{L}_n, \quad \mathcal{L}_1 = w' - w^2. \quad (26)$

Example 8. (*n* = 2) $\left(w'''' - 10w^2w'' - 10w{w'}^2 + 6w^5\right) + t_1\left(w'' - 2w^3\right) = zw + \alpha_2$

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