Coverings and multivector pseudosymmetries of differential equations

Vladimir Chetverikov

Bauman Moscow State Technical University

Posing the problem: coverings over equations

$$\mathcal{E}: F(z, v, v_{\sigma}) = 0$$

$$v = (v^1, \dots, v^m), \qquad z = (z_1, \dots, z_n),$$
$$v_{\sigma} = \left(\frac{\partial^{|\sigma|} v^1}{\partial z_1^{i_1} \dots \partial z_n^{i_n}}, \dots, \frac{\partial^{|\sigma|} v^m}{\partial z_1^{i_1} \dots \partial z_n^{i_n}}\right), \quad \sigma = (i_1, \dots, i_n).$$

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Posing the problem: coverings over equations

$$\tilde{\mathcal{E}}: \begin{cases} F(z, v, v_{\sigma}) = 0, \\ \frac{\partial w}{\partial z_i} = W_i(z, w, v, v_{\sigma}), & i = \overline{1, n}, \\ w = (w^1, \dots, w^q) \end{cases}$$

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The system $\tilde{\mathcal{E}}$ covers the system \mathcal{E} , ν is a covering over \mathcal{E} .

Posing the problem: coverings from equations

$$\mathcal{S}: G(x, u, u_{\sigma}) = 0, \qquad x = (x_1, \dots, x_n), \quad u = (u^1, \dots, u^{m_0})$$

 $(z, v, w) = \phi(x, u, u_{\sigma}) \downarrow \qquad \uparrow (x, u) = \phi^{-1}(z, v, w, v_{\sigma})$

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 ϕ is a *C*-transformation, the system *S* covers the system $\mathcal{E}, \nu \circ \phi$ is a covering from *S*, the fields $\partial/\partial w^i$ constitute a pseudosymmetry of *S*, the system $\tilde{\mathcal{E}}$ is a decomposable form of the covering.

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Example 1: the Laplace equation

$$\mathcal{S}: u_{(2,0)} + u_{(0,2)} = 0, \quad u_{(2,0)}^2 + u_{(1,1)}^2 \neq 0,$$

the C-transformation ϕ :

$$z_1 = u_{(1,0)}, \quad z_2 = u_{(0,1)}, \quad v = u, \quad w^1 = x_1, \quad w^2 = x_2,$$

the inverse transformation ϕ^{-1} : $x_1 = w^1$, $x_2 = w^2$, u = v, the decomposable form:

$$\begin{split} \tilde{\mathcal{E}}: \quad v_{(2,0)} + v_{(0,2)} &= 0, \quad w_{(1,0)}^1 = \frac{z_1 v_{(1,0)} - z_2 v_{(0,1)}}{z_1^2 + z_2^2} = -w_{(0,1)}^2, \\ w_{(0,1)}^1 &= \frac{z_1 v_{(0,1)} + z_2 v_{(1,0)}}{z_1^2 + z_2^2} = w_{(1,0)}^2, \end{split}$$

the pseudosymmetry: $(\partial/\partial x_1, \partial/\partial x_2)$.

Outline

Pseudosymmetries of differential equations

Description of coverings from equations

Multivector pseudosymmetries

Examples

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- V. V. Sokolov.

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The spaces of infinite jets

Consider an infinite-dimensional space J^{∞} with coordinates

$$x_i, u^j, u^j_{\sigma}, \qquad i = \overline{1, n}, \ j = \overline{1, m}, \ |\sigma| \ge 0.$$

The total derivative with respect to x_i on J^{∞} :

$$D_i = \frac{\partial}{\partial x_i} + \sum_{j=1}^m \sum_{|\sigma| \ge 0} u^j_{\sigma+1_i} \frac{\partial}{\partial u^j_{\sigma}}.$$

The infinite prolongation of system $S : G_{\alpha}(x, u, u_{\sigma}) = 0, \alpha = \overline{1, r}$, is given by

$$\mathcal{S}^{\infty}: \quad D_{\sigma}G_{\alpha} = 0, \qquad |\sigma| \ge 0, \quad \alpha = \overline{1, r},$$

where $D_{\sigma} = D_1^{i_1} \circ \ldots \circ D_n^{i_n}$ for $\sigma = (i_1, \ldots, i_n), S^{\infty} \subset J^{\infty}$.

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The Cartan distribution

The Cartan distribution on S^{∞} :

$$CD(S) = \operatorname{span}_{C^{\infty}(S^{\infty})} \{ D_1 |_{S^{\infty}}, \dots, D_n |_{S^{\infty}} \}.$$

The Cartan forms on \mathcal{S}^{∞} :

$$d_{\mathcal{C}}g = \mathrm{d}g - \sum_{i=1}^{n} D_{i}(g)\mathrm{d}x_{i}, \quad g \in C^{\infty}(\mathcal{S}^{\infty}).$$
$$\mathcal{C}\Lambda^{1}(\mathcal{S}) = \mathrm{span}_{C^{\infty}(\mathcal{S}^{\infty})}\{d_{\mathcal{C}}g : g \in C^{\infty}(\mathcal{S}^{\infty})\}.$$
$$\mathcal{C}^{k}\Lambda^{k}(\mathcal{S}) = \underbrace{\mathcal{C}\Lambda^{1}(\mathcal{S}) \wedge \cdots \wedge \mathcal{C}\Lambda^{1}(\mathcal{S})}_{k \text{ times}}.$$
$$w \in \mathcal{C}^{k}\Lambda^{k}(\mathcal{S}) \iff D_{i}|_{k} |w| = 0, \quad i = \overline{1, n}.$$

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Mappings of diffieties

A mapping $\phi : \mathcal{S}^{\infty} \longrightarrow \mathcal{E}^{\infty}$ is *C*-transformation if

- ϕ is smooth, i.e., $\phi^*(C^{\infty}(\mathcal{E}^{\infty})) \subset C^{\infty}(\mathcal{S}^{\infty}), \phi^*(g) = g \circ \phi;$
- there exists a smooth inverse mapping;
- $\phi_*(\mathcal{CD}(\mathcal{S})) = \mathcal{CD}(\mathcal{E}).$

A smooth mapping $\nu : S^{\infty} \longrightarrow \mathcal{E}^{\infty}$ is a covering if

- the tangent mapping ν_{*,θ} is a vector space epimorphism for any θ ∈ S[∞];
- $\nu_*(\mathcal{CD}(S)) = \mathcal{CD}(\mathcal{E});$
- the dim ker $\nu_{*,\theta}$ is the same for any $\theta \in S^{\infty}$ (rank ν).

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Conditions for fibers of coverings

$$F_{\beta}(z, v, \dots, v_{\sigma}, \dots) = 0, \quad \beta = \overline{1, r_1},$$

$$\frac{\partial w^j}{\partial z_i} = W_i^j(z, w, v, \dots, v_{\sigma}, \dots), \quad i = \overline{1, n}, \ j = \overline{1, q},$$
(1)
(2)

The total derivative with respect to x_i for system (1)–(2) has the form

$$D_i = \check{D}_i + \sum_{j=1}^q W_i^j(x, w, u, \dots, u_\sigma, \dots) \frac{\partial}{\partial w^j},$$

where \check{D}_i is the total derivative w.r.t. x_i for (1).

$$[X, D_i] = A_i X, \qquad i = \overline{1, n}, \tag{3}$$

where $X = \left(\partial/\partial w^1, \dots, \partial/\partial w^q\right)^{\mathrm{T}}, A_i = \left(\partial W_i^j/\partial w^s\right)_{\mathrm{T}}$

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The invariant version of (3):

$$[X, D_i] = A_i X + B_i D, \qquad i = \overline{1, n}, \tag{4}$$

where $D = (D_1, ..., D_n)^T$. A column *X* of vector fields satisfying (4) is a pseudosymmetry of the system *S*. The relations (4) and the Jacobi identity imply:

$$D_{j}(A_{i}) - D_{i}(A_{j}) + A_{j}A_{i} - A_{i}A_{j} = 0, \qquad \forall i, j, \qquad (5) D_{j}(B_{i}) - D_{i}(B_{j}) + A_{j}B_{i} - A_{i}B_{j} = 0, \qquad \forall i, j.$$

Equations (5) are the Maurer–Cartan equations.

$$(5) \implies [D_i + A_i, D_j + A_j] = 0.$$

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Pseudosymmetries of infinite jets spaces

Theorem 1. A *q*-column *X* is a pseudosymmetry of J^{∞} iff

$$X = \Im_{\varphi,A} + MD,$$

where φ is a $q \times m$ matrix of functions on J^{∞} , $A = (A_1, \ldots, A_n)$ is a tuple of $q \times q$ matrices satisfying (5), M is a $q \times n$ matrix of arbitrary functions on J^{∞} , the term $\mathfrak{D}_{\varphi,A}$ has the form

$$\Im_{\varphi,A} = \sum_{\sigma} (D+A)^{\sigma}(\varphi) \frac{\partial}{\partial u_{\sigma}},$$

where $(D+A)^{\sigma} = (D_1 + A_1)^{i_1} \circ \ldots \circ (D_n + A_n)^{i_n}$, $\sigma = (i_1, \ldots, i_n)$, is a $(q \times q)$ matrix differential operator, $\frac{\partial}{\partial u_{\sigma}} = \left(\frac{\partial}{\partial u_{\sigma}^1}, \ldots, \frac{\partial}{\partial u_{\sigma}^m}\right)^{\mathrm{T}}$. Defining equations for pseudosymmetries

Theorem 2. Let

$$\mathcal{S}:G(x,u,u_{\sigma})=0$$

be a formally integrable system. Then a column *X* is a pseudosymmetry of *S* iff *X* is the restriction of $\Im_{\varphi,A} + MD$ to S^{∞} , where the matrix $\varphi = (\varphi_{ij})$ satisfies

$$\sum_{j,\sigma} \frac{\partial G}{\partial u_{\sigma}^{j}} (D+A)^{\sigma} (\varphi^{j}) \Big|_{\mathcal{S}^{\infty}} = 0, \quad \varphi^{j} = (\varphi_{1j}, \dots, \varphi_{qj})^{\mathrm{T}}.$$

The matrix φ is the generating matrix of the pseudosymmetry, and $\Im_{\varphi,A}$ is the evolution pseudosymmetry.

Defining equations for pseudosymmetries

Theorem 2. Let

$$\mathcal{S}:G(x,u,u_{\sigma})=0$$

be a formally integrable system. Then a column X is a pseudosymmetry of S iff X is the restriction of $\Im_{\varphi,A} + MD$ to S^{∞} , where the matrix $\varphi = (\varphi_{ij})$ satisfies

$$\sum_{j,\sigma} \frac{\partial G}{\partial u_{\sigma}^{j}} \left(D + A \right)^{\sigma} (\varphi^{j}) \Big|_{\mathcal{S}^{\infty}} = 0, \quad \varphi^{j} = (\varphi_{1j}, \dots, \varphi_{qj})^{\mathrm{T}}.$$

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Integrable pseudosymmetries

A column X of vector fields X_1, \ldots, X_q on S^{∞} is an integrable pseudosymmetry of the system S if

• $X = (X_1, \ldots, X_q)^T$ is a pseudosymmetry of S;

• X_1, \ldots, X_q generate an involutive distribution of dimension q;

- the finiteness condition: there exists a ring $\mathcal K$ such that
 - (a) $\mathcal{F}_0(\mathcal{S}) \subset \mathcal{K} \subset \mathcal{F}_l(\mathcal{S})$ for some $l \ge 0$;
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Here $\mathcal{F}_i(\mathcal{S})$ consists of smooth functions of $t, u, \ldots, u_{\sigma}, |\sigma| \leq i$. Example 2.

The fields $\partial/\partial w^1, \ldots, \partial/\partial w^q$ form an integrable pseudosymmetry of system (1)–(2), \mathcal{K} is the ring of smooth functions of

 $t, w, v, \ldots, v_{\sigma}, |\sigma| \leq s$, where *s* is the order of system (2).

A column X of vector fields X_1, \ldots, X_q on S^{∞} is an integrable pseudosymmetry of the system S if

- $X = (X_1, \ldots, X_q)^T$ is a pseudosymmetry of S;
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An integrable pseudosymmetry $X = (X_1, ..., X_q)^T$ determines a covering if the fibers of the covering coincide with the maximal integral manifolds of the distribution generated by $X_1, ..., X_q$.

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Description of coverings

Let S be a formally integrable system, and let $X = (X_1, \ldots, X_q)^T$ be an integrable pseudosymmetry of S. A point $\theta \in S^\infty$ is a regular point of X if

- the fields $X_1, \ldots, X_q, D_1, \ldots, D_n$ are linearly independent at θ ;
- the subspaces {df|_{θ'} ∈ T^{*}_{θ'} : f ∈ K} and {df|_{θ'} ∈ T^{*}_{θ'} : f ∈ D^lK} have constant dimension in some neighborhood of θ, where D^lK is the ring generated by the functions D_σ(f), f ∈ K, |σ| ≤ l, F₀(S) ⊂ K ⊂ F_l(S).

Theorem 3.

- If X is an integrable pseudosymmetry of a system S and θ ∈ S[∞] is a regular point of X, then X determines a covering from a neighborhood of θ.
- For any covering from a system *S* there exists an integrable pseudosymmetry of *S* determining this covering.

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Let $X = (X_1, \ldots, X_q)^T$ be an integrable pseudosymmetry of a system S and let θ be a regular point of X.

Step 1. Find an integer *s* such that the coefficients of matrices A_1, \ldots, A_n and B_1, \ldots, B_n belong to $D^s \mathcal{K}$. Step 2. Find functions g_1, \ldots, g_p such that

1. $g_1,\ldots,g_p\in D^s\mathcal{K};$

- 2. g_1, \ldots, g_p are common first integrals of X_1, \ldots, X_q ;
- 3. dg_1, \ldots, dg_p are linearly independent at each point;
- 4. The set $\{g_1, \ldots, g_p\}$ is a maximal set satisfying 1-3.

Step 3. Find functions $w_1, \ldots, w_q \in \mathcal{F}(S)$ such that the matrix $(X_i(w_j)(\theta))$ is nonsingular. Step 4. Choose new independent variables x_1, \ldots, x_n among the functions g_1, \ldots, g_p such that the matrix $(D_i(x_j)(\theta))$ is nonsingular. By *u* denote the set of the functions g_1, \ldots, g_p that are not $x_1, \ldots, x_n \in \mathbb{R}$.

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Integrable pseudosymmetries

Theorem 4. Suppose a column $X = (X_1, \ldots, X_q)^T$ is an integrable pseudosymmetry of a system S and N is a nonsingular $q \times q$ matrix of functions on S^{∞} . Then

- the product *NX* is also an integrable pseudosymmetry of the system *S*;
- the integrable pseudosymmetries X and NX determine the same covering;

Methods for solving equations for pseudosymmetries are similar to methods for solving equations for higher symmetries. To find integrable pseudosymmetries, these methods should be combined with using Theorem 4 and the finiteness condition.

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- if matrices A_i and B_i correspond to X (see (4)), then the matrices $NA_iN^{-1} D_i(N)N^{-1}$ and NB_i correspond to NX.

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Commutation of pseudosymmetries

Let *X* and *Y* be two pseudosymmetries of a system S. Denote by $[X \cup Y]$ the set of generators of the involutive closure of the union of the sets *X* and *Y* of vector fields.

Theorem 5.

If *X* and *Y* are two pseudosymmetries of a system *S* and the set $[X \cup Y]$ is finite, then this set also forms a pseudosymmetry of the system *S*.

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Multivector fields and their Schouten bracket

The *r*-th exterior degree of the module $\mathcal{D}(\mathcal{S}^{\infty})$ of vector fields on \mathcal{S}^{∞} :

$$\bigwedge^{r} \mathcal{D}(\mathcal{S}^{\infty}) = \operatorname{span}_{\mathcal{F}(\mathcal{S})} \{ X_{1} \wedge \cdots \wedge X_{r} : X_{1}, \dots, X_{r} \in \mathcal{D}(\mathcal{S}^{\infty}) \}.$$

Elements of the module $\bigwedge^* \mathcal{D}(\mathcal{S}^\infty) = \bigoplus_{r>0} \bigwedge^r \mathcal{D}(\mathcal{S}^\infty)$ are multivector fields on \mathcal{S}^∞ . An element $X \in \bigwedge^r \mathcal{D}(\mathcal{S}^\infty)$ is a homogeneous element of degree |X| = r. The Schouten bracket of $X = \bigwedge^s X_s$ and $Y = \bigwedge^k Y_s$ is

$$[X,Y] = \sum_{\alpha=1}^{s} \sum_{\beta=1}^{k} (-1)^{\alpha+\beta} [X_{\alpha}, Y_{\beta}] \wedge X_{1} \wedge \dots \wedge \widehat{X_{\alpha}} \wedge \dots$$

This bracket is extended to $igwedge^*\mathcal{D}(\mathcal{S}^\infty)$ by linearity.

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This bracket is extended to $\bigwedge^* \mathcal{D}(\mathcal{S}^\infty)$ by linearity.

Properties of the Schouten bracket

Theorem 6. The following identities are valid for homogeneous elements X, Y, Z:

- [*X*, *Y*] = −(−1)^{(|*X*|−1)(|*Y*|−1)}[*Y*, *X*] (antisymmetry of Schouten bracket);
- $(-1)^{(|X|-1)(|Z|-1)}[X, [Y, Z]] + (-1)^{(|Y|-1)(|X|-1)}[Y, [Z, X]] + (-1)^{(|Z|-1)(|Y|-1)}[Z, [X, Y]] = 0$ (Jacobi identity for Schouten bracket);
- $[X \wedge Y, Z] = (-1)^{|Y|(|Z|-1)}[X, Z] \wedge Y + X \wedge [Y, Z]$ (Poisson identity).

The Schouten bracket makes the module $\bigwedge^* \mathcal{D}(\mathcal{S}^\infty)$ into a Lie superalgebra if the degree of a homogeneous element $X \in \bigwedge^r \mathcal{D}(\mathcal{S}^\infty)$ is taken to be the integer r - 1.

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Theorem 7. If $(X_1, \ldots, X_q)^T$ is a pseudosymmetry of system S, then $X = X_1 \land \cdots \land X_q$ satisfies the equalities

$$[X, D_i] = a_{i,X}X + \sum_{j=1}^n Z_{ij,X} \wedge D_j, \qquad i = \overline{1, n}, \tag{6}$$

for some $a_{i,X} \in C^{\infty}(\mathcal{S}^{\infty})$ and $Z_{ij,X} \in \bigwedge^{q-1} \mathcal{D}(\mathcal{S}^{\infty})$.

A multivector field X on S^{∞} satisfying the equality (6) is a multivector pseudosymmetry of system S.

Theorem 8. If X and Y are multivector pseudosymmetries of a system, then $X \wedge Y$ is also a multivector pseudosymmetry of this system.

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Multivector pseudosymmetries

The interior product of a multivector field $X = X_1 \land \cdots \land X_r$ and a *p*-form $\omega, p \ge r$:

$$i_X(\omega) = (i_{X_r} \circ \cdots \circ i_{X_1})(\omega),$$

where i_{X_i} is the interior product of the vector field X_i . Theorem 9.

$$i_X \circ Z - Z \circ i_X = i_{[X,Z]}, \quad X \in \bigwedge^* \mathcal{D}(\mathcal{S}^\infty), \quad Z \in \mathcal{D}(\mathcal{S}^\infty).$$

Theorem 10. If X is a multivector pseudosymmetry of a system S, $q = |X|, \omega \in C^q \Lambda^q(S^\infty)$, then

$$i_X(D_l\omega) = (D_l + a_{l,X})(i_X\omega).$$

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A scalar C-differential operator of $C^q \Lambda^q(\mathcal{S}^\infty)$ to $C^q \Lambda^q(\mathcal{S}^\infty)$:

$$\omega \mapsto \sum_{|\sigma| \le k} a_{\sigma} D_{\sigma}(\omega), \quad a_{\sigma} \in \mathcal{F}(\mathcal{S}).$$

The set of all scalar C-differential operators is a noncommutative ring with respect to the composition operation, $C^q \Lambda^q(S^\infty)$ is a left module over this ring.

Theorem 10 \implies a multivector pseudosymmetry *X* of a system *S* is uniquely determined by its values at generators of the module $C^q \Lambda^q(S^\infty), q = |X|.$

Example 3. For the Laplace equation $u_{x_1x_1} + u_{x_2x_2} = 0$ generators of the module $C^2 \Lambda^2(S^{\infty})$ are

$$\omega_1^k = -d_{\mathcal{C}}u_{(k,0)} \wedge d_{\mathcal{C}}u_{(k,1)}, \quad \omega_2^k = d_{\mathcal{C}}u_{(k,0)} \wedge d_{\mathcal{C}}u_{(k+1,0)}, \quad k \ge 0.$$

They are related by the equations

$$D_2(\omega_1^k) = D_1(\omega_2^k), \quad \Delta^k \left(D_1(\omega_1^0) + D_2(\omega_2^0) \right) = c_k \left(D_1(\omega_1^k) + D_2(\omega_2^k) \right),$$

where $\{c_k\}$ is some sequence of integers, $\Delta = D_1^2 + D_2^2$. If *X* is the multivector pseudosymmetry such that

$$i_X(\omega_1^0) = x_2, \quad i_X(\omega_2^0) = x_1, \quad i_X(\omega_i^k) = 0, \quad a_{i,X} = 0, \quad Z_{ij,X} = 0$$

for any k > 0, i, j = 1, 2. Then

$$X = x_1 \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(1,0)}} - x_2 \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(0,1)}} + \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(2,0)}}.$$

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Constructing coverings determined by a multivector pseudosymmetry

The Lie derivative of $g \in \mathcal{F}(S)$ along $X = X_1 \wedge \cdots \wedge X_q$:

$$X(g) = \sum_{i=1}^{q} (-1)^{q-i} X_i(g) X_1 \wedge \dots \wedge \widehat{X}_i \wedge \dots \wedge X_q.$$

A function $g \in \mathcal{F}(\mathcal{S})$ is a first integral of *X* if $X(g) \equiv 0$.

 $(X_1 \wedge \cdots \wedge X_q)(g) \equiv 0 \quad \Longleftrightarrow \quad X_i(g) \equiv 0, \quad i = \overline{1, q}.$

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$$(X_1 \wedge \cdots \wedge X_q)(g) \equiv 0 \quad \Longleftrightarrow \quad X_i(g) \equiv 0, \quad i = \overline{1, q}.$$

Example 4. For

$$X = x_1 \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(1,0)}} - x_2 \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(0,1)}} + \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(2,0)}}$$

the corresponding decomposable form is

$$(z_2 - z_1)(z_2v_{(1,0)}^1 - z_1v_{(0,1)}^1) = (z_2^2 + z_1^2)v^2 - (z_2 + z_1)v^1,$$

$$w_{(1,0)}^1 = \frac{v^1 - z_1w^2}{z_2 - z_1}, \qquad w_{(0,1)}^1 = \frac{v^1 - z_2w^2}{z_1 - z_2},$$

the corresponding change of variables is

$$z_1 = x_1, \quad z_2 = x_2, \quad w^1 = u, \quad w^2 = u_{(1,0)} + u_{(0,1)},$$

 $v^1 = x_2 u_{(1,0)} + x_1 u_{(0,1)}, \quad v^2 = u_{(1,0)} + u_{(0,1)} + (x_2 - x_1) u_{(2,0)}.$

Example 5. In the situation of Example 3, but when $i_X(\omega_2^1) = -1$ and $i_X(\omega_i^k) = 0$ for other k, i, we have

$$X = \frac{\partial}{\partial u} \wedge \frac{\partial}{\partial u_{(3,0)}} - \frac{\partial}{\partial u_{(1,0)}} \wedge \frac{\partial}{\partial u_{(2,0)}} + \frac{\partial}{\partial u_{(0,1)}} \wedge \frac{\partial}{\partial u_{(1,1)}}.$$

The decomposable form for X is

$$v_{(1,0)}^1 = v_{(0,1)}^2, \quad v_{(0,1)}^1 = -v_{(1,0)}^2, \quad w_{(2,1)} = v^1, \quad w_{(4,0)} = v^2.$$

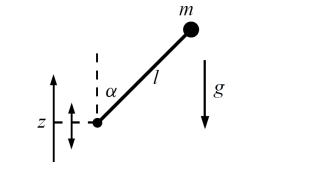
The change of variables is

$$z_1 = x_1, \quad z_2 = x_2, \quad v^1 = u_{(2,1)}, \quad v^2 = u_{(4,0)}, \quad w = u.$$

Thus the dimension of a covering can be greater than the degree of the corresponding multivector pseudosymmetry.

医马克氏 医黄疸

Example 6: the Kapitsa pendulum



$$\dot{\alpha} = p + \frac{u}{l}\sin\alpha, \quad \dot{z} = u, \quad \dot{p} = \left(\frac{g}{l} - \frac{u^2}{l^2}\cos\alpha\right)\sin\alpha - \frac{u}{l}p\cos\alpha.$$

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This system is equivalent to the system

$$\ddot{\mathbf{y}} = a(\alpha) \dot{\alpha}^2, \quad a(\alpha) = l \frac{\cos \alpha}{\sin^2 \alpha},$$

where

$$y = z + g\frac{t^2}{2} - l \int \frac{d\alpha}{\sin\alpha}.$$

Let us compute all the one-dimensional integrable evolution pseudosymmetries $\mathfrak{D}_{\varphi,A}$ of this system.

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Denote by \mathcal{F}_s is the ring of smooth functions depending on $t, y, \alpha, \dot{y}, \dot{\alpha}, \dots, \alpha^{(s)}$ for $s \ge 0$. The defining equation for pseudosymmetries is

$$(D_t + A)^2(\varphi_y) = 2a \dot{\alpha} (D_t + A)(\varphi_\alpha) + a' \dot{\alpha}^2 \varphi_\alpha,$$

where $\varphi = (\varphi_y, \varphi_\alpha)^T$ is a generating matrix and $A, \varphi_y, \varphi_\alpha \in \mathcal{F}_s$ for some $s \ge 0$.

The three cases are possible:

I) $\varphi_y \neq 0, \ \varphi_\alpha \neq 0;$ II) $\varphi_y \neq 0, \ \varphi_\alpha \equiv 0;$ III) $\varphi_y \neq 0, \ \varphi_\alpha \equiv 0;$ III) $\varphi_y \equiv 0, \ \varphi_\alpha \neq 0.$

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In case I) the column $(1, \psi = \varphi_{\alpha}/\varphi_y)^T$ is also a generating matrix of an integrable pseudosymmetry and

$$D_{t}(A) + A^{2} = 2a(\alpha) \dot{\alpha} (D_{t} + A)(\psi) + a'(\alpha) \dot{\alpha}^{2} \psi.$$
(7)
$$\psi \in \mathcal{F}_{s} \setminus \mathcal{F}_{s-1} \implies A \in \mathcal{F}_{s} \setminus \mathcal{F}_{s-1}.$$

$$\partial_{\psi,A} = \dots + \sum_{k \ge 0} (D_{J} + A)^{k}(\psi) \frac{\partial}{\partial \alpha^{(k)}} \implies$$

$$\partial_{\psi,A} \colon \mathcal{F}_{k} \to \mathcal{F}_{s+k} \implies s = 0 \text{ (from the finiteness condition)}.$$

Thus, ψ and A are functions of t, y, α, \dot{y} only.

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$$\psi \in \mathcal{F}_{s} \setminus \mathcal{F}_{s-1} \implies A \in \mathcal{F}_{s} \setminus \mathcal{F}_{s-1}.$$

$$\Im_{\psi,A} = \dots + \sum_{k \ge 0} (D_{J} + A)^{k}(\psi) \frac{\partial}{\partial \alpha^{(k)}} \implies$$

$$\Im_{\psi,A}: \mathcal{F}_{k} \to \mathcal{F}_{s+k} \implies s = 0 \text{ (from the finiteness condition).}$$

Thus, ψ and A are functions of t, y, α, \dot{y} only.

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Finding successively the coefficients at $\dot{\alpha}^3$, $\dot{\alpha}^2$, \dot{y}^2 , $\dot{y}\dot{\alpha}$, $\dot{\alpha}$, and \dot{y} in (7), we get

$$\psi = \frac{b+c_1}{2\sqrt{a}(y+c_2)}, \ A = \frac{\dot{y}}{y+c_2}, \ b = \int \sqrt{a}d\alpha,$$

where c_1, c_2 are arbitrary constants.

Similarly, in the cases II) and III) we get the following pseudosymmetries:

II) $\varphi_y = 1, \ \varphi_\alpha = 0, \ A = \frac{1}{t+c_3}, \ c_3 = \text{const, on}$ $\varphi_y = 1, \ \varphi_\alpha = 0, \ A = 0;$ III) $\varphi_y = 0, \ \varphi_\alpha = 1, \ A = -\frac{a' \dot{\alpha}}{2a}.$

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These pseudosymmetries define the following decomposable forms:

1)
$$\begin{cases} \dot{v}_{2} = (\dot{v}_{1}^{2} + \frac{1}{2}v_{1}v_{2})^{2} - v_{2}^{2} & v_{1} = \frac{\beta}{\sqrt{y+c_{2}}}, \quad v_{2} = \frac{\dot{y}}{y+c_{2}}, \\ \dot{w} = v_{2}(w+c_{2}), & w = y; \end{cases}$$
2)
$$\begin{cases} \dot{v}_{1} = -(t+c_{3})a(v_{2})\dot{v}_{2}^{2} & v_{1} = y - \dot{y}(t+c_{3}), \quad v_{2} = \alpha, \\ \dot{w} = \frac{w-v_{1}}{t+c_{3}}, & w = y; \end{cases}$$

$$\begin{cases} \dot{v}_{1} = a(v_{2})\dot{v}_{2}^{2} & v_{1} = \dot{y}, \quad v_{2} = \alpha, \\ \dot{w} = v_{1}, & w = y; \end{cases}$$
3)
$$\dot{w} = \sqrt{\frac{\ddot{v}}{a(w)}}, & v = y, \\ w = \alpha. \end{cases}$$

Commutation of the obtained pseudosymmetries

Let \mathfrak{D}_1 and \mathfrak{D}_3 be the pseudosymmetries corresponding to cases (1) and (3), and let \mathfrak{D}_2 and \mathfrak{D}_4 be the pseudosymmetry corresponding to case (2). Then

$$\begin{split} [\Im_1, \Im_2] &= \frac{1}{y + c_2} (\Im_1 - \Im_2), \quad [\Im_1, \Im_3] = \frac{h(\alpha)}{y + c_2} \Im_3, \\ [\Im_1, \Im_4] &= \frac{1}{y + c_2} (\Im_1 - \Im_4), \quad [\Im_2, \Im_3] = 0, \quad [\Im_2, \Im_4] = 0, \\ [\Im_3, \Im_4] &= 0, \quad h(\alpha) = \frac{(b + c_1)a'}{4a^{3/2}} - \frac{1}{2}. \end{split}$$

Thus, any nonempty subset of the tuple $\{\Im_1, \Im_2, \Im_3, \Im_4\}$ generates an integrable pseudosymmetry of the system.

We obtain 6 two-dimensional, 4 three-dimensional, and 1 four-dimensional integrable pseudosymmetries and 11 coverings from the system. The decomposable forms corresponding to the pseudosymmetries $(\mathfrak{D}_1, \mathfrak{D}_3), (\mathfrak{D}_1, \mathfrak{D}_2, \mathfrak{D}_3),$ and $(\mathfrak{D}_1, \mathfrak{D}_2, \mathfrak{D}_3, \mathfrak{D}_4)$ are:

$$\begin{cases} \dot{w}_{1} = v(w_{1} + c_{2}) \\ \dot{w}_{2} = \sqrt{\frac{(\dot{v} + v^{2})(w_{1} + c_{2})}{a(w_{2})}}, & v = \frac{\dot{y}}{y + c_{2}}, w_{1} = y, w_{2} = \alpha; \\ \\ \dot{w}_{1} = \frac{\dot{w}_{1}(t + c_{3}) - w_{1} - c_{2}}{v} \\ \dot{w}_{2} = \sqrt{\frac{\dot{w}_{1}(t + c_{3}) - w_{1} - c_{2}}{va(w_{2})}}, & v = \frac{\dot{y}(t + c_{3}) - y - c_{2}}{a(\alpha) \dot{\alpha}^{2}}, w_{1} = y, w_{2} = \alpha; \\ \\ \dot{w}_{2} = \sqrt{\frac{\dot{w}_{1}(t + c_{3}) - w_{1} - c_{2}}{va(w_{2})}}, & v = \frac{\ddot{y}(t + c_{3}) - y - c_{2}}{a(\alpha) \dot{\alpha}^{2}}, w_{1} = y, w_{2} = \alpha; \\ \\ \\ \ddot{w}_{1} = v\ddot{w}_{1} \\ \dot{w}_{2} = \sqrt{\frac{\ddot{w}_{1}}{a(w_{2})}}, & v = \frac{\ddot{y}}{\ddot{y}}, w_{1} = y, w_{2} = \alpha. \end{cases}$$

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THANK YOU FOR YOUR ATTENTION

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