

Projective-geometric aspects of WDVV equations

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Local and Nonlocal Geometry of PDEs
and Integrability – Oct. 2018

Contents

- ▶ WDVV equations and their bi-Hamiltonian structure.
- ▶ Projective-geometric interpretation.
- ▶ Open questions.

Witten–Dijkgraaf–Verlinde–Verlinde equations

The problem: in \mathbb{R}^N find a function $F = F(t^1, \dots, t^N)$ such that

1. $\frac{\partial^3 F}{\partial t^1 \partial t^\alpha \partial t^\beta} = \eta_{\alpha\beta}$ constant symmetric nondegenerate matrix
2. $c_{\alpha\beta}^\gamma = \eta^{\gamma\epsilon} \frac{\partial^3 F}{\partial t^\epsilon \partial t^\alpha \partial t^\beta}$ structure constants of an associative algebra
3. $F(c^{d_1} t^1, \dots, c^{d_N} t^N) = c^{d_F} F(t^1, \dots, t^N)$ quasihomogeneity ($d_1 = 1$)

If e_1, \dots, e_N is the basis of \mathbb{R}^N then the algebra operation is

$$e_\alpha \cdot e_\beta = c_{\alpha\beta}^\gamma(\mathbf{t}) e_\gamma \quad \text{with unity } e_1$$

WDVV equations of associativity

$$\eta^{\mu\lambda} \frac{\partial^3 F}{\partial t^\lambda \partial t^\alpha \partial t^\beta} \frac{\partial^3 F}{\partial t^\nu \partial t^\mu \partial t^\gamma} = \eta^{\mu\lambda} \frac{\partial^3 F}{\partial t^\nu \partial t^\alpha \partial t^\mu} \frac{\partial^3 F}{\partial t^\lambda \partial t^\beta \partial t^\gamma} \quad (\text{WDVV})$$

Why study WDVV?

1. Solutions yield Gromov–Witten invariants
2. Solutions correspond to integrable hierarchies
3. Applications to Quantum Field Theory (?)

WDVV equations in detail

$d_F \neq 3$: By linear transformations preserving e_1 :

$$\eta_{\alpha\beta} = \delta_{\alpha+\beta, N+1} = \begin{pmatrix} 0 & & 1 \\ & \ddots & \\ 1 & & 0 \end{pmatrix}$$

$$F = \frac{1}{2}(t^1)^2 t^N + \frac{1}{2} t^1 \sum_{\alpha=2}^{N-1} t^\alpha t^{N-\alpha+1} + f(t^2, \dots, t^N);$$

$d_F = 3$: By linear transformations preserving e_1 :

$$\eta_{\alpha\beta} = \delta_{\alpha\beta}$$

$$F = \frac{1}{6}(t^1)^3 + \frac{1}{2} t^1 \sum_{\alpha=2}^N (t^\alpha)^2 + f(t^2, \dots, t^N);$$

Main example: WDVV in the case $N = 3$

If $N = 3$ we have a single equation on $f = f(t^2, t^3) = f(x, t)$.

Two cases:

- ▶ $\eta_{\alpha\beta} = \delta_{\alpha+\beta, N+1}$, then

$$f_{ttt} = f_{xxt}^2 - f_{xxx}f_{xtt}$$

- ▶ $\eta_{\alpha\beta} = \delta_{\alpha\beta}$, then

$$f_{ttt} = \frac{f_{xxt}^2 - f_{xxx}f_{xtt} + f_{xtt}^2 - 1}{f_{xxt}}$$

WDVV equations as hydrodynamic systems

Construction by O. Mokhov (1995). Let us introduce coordinates

$$a = f_{xxx}, \quad b = f_{xxt}, \quad c = f_{xtt}.$$

Then the compatibility conditions in the two cases are

$$\left\{ \begin{array}{l} a_t = b_x, \\ b_t = c_x, \\ c_t = (b^2 - ac)_x \end{array} \right. \quad \text{and} \quad \left\{ \begin{array}{l} a_t = b_x, \\ b_t = c_x, \\ c_t = \left(\frac{-ac + b^2 + c^2 - 1}{b} \right)_x \end{array} \right.$$

The system on the left is bi-Hamiltonian (Ferapontov, Galvao, Mokhov, Nutku CMP'98) by a third-order and a first-order Hamiltonian operator of Dubrovin–Novikov type. Further study by Kersten, Krasil'shchik, Verbovetsky, V. TMP'12. **What about the system on the right?**

Digression: third-order Hamiltonian operators

Homogeneous, or Dubrovin–Novikov Hamiltonian operator ($\mathbf{u} = (u^1, \dots, u^n)$):

$$\begin{aligned} A^{ij} = & g^{ij}(\mathbf{u}) \partial_x^3 + b_k^{ij}(\mathbf{u}) u_x^k \partial_x^2 \\ & + [c_k^{ij}(\mathbf{u}) u_{xx}^k + c_{km}^{ij}(\mathbf{u}) u_x^k u_x^m] \partial_x \\ & + d_k^{ij}(\mathbf{u}) u_{xxx}^k + d_{km}^{ij}(\mathbf{u}) u_x^k u_{xx}^m + d_{kmn}^{ij}(\mathbf{u}) u_x^k u_x^m u_x^n. \end{aligned}$$

They admit the **canonical form** (Potemin '89,'92; Doyle '92)

$$A^{ij} = \partial_x \left(g^{ij} \partial_x + c_k^{ij} u_x^k \right) \partial_x$$

They correspond to **quadratic line complexes**, and are projectively invariant (Ferapontov, Pavlov, V., JGP '14, IMRN '16).

Digression: Plücker's line geometry

Two points $U, V \in \mathbb{P}^n(\mathbb{C})$, $U = [u^0, \dots, u^n]$, $V = [v^0, \dots, v^n]$ define a line with coordinates $p_{ij} = \det \begin{vmatrix} u^i & u^j \\ v^i & v^j \end{vmatrix}$ inside a projective space: $\mathbb{P}^n(\mathbb{C}) \hookrightarrow \mathbb{P}^N(\mathbb{C})$ (Plücker's embedding).

A **quadratic line complex** is a variety

$$X^T Q X = f_{ij,hk} p^{ij} p^{hk} = 0$$

Take $V = U + dU$: then $p^{ij} = u^i du^j - u^j du^i$ and we have a **Monge metric**

$$g_{ab} du^a du^b = f_{ij,hk} (u^i du^j - u^j du^i)(u^h du^k - u^k du^h)$$

We set $u^0 = 1$, $du^0 = 0$.

Digression: hydrodynamic systems of cons. laws

Theorem: (Fereapontov, Pavlov, V. LMP '17) Let A be a third-order homogeneous Hamiltonian operator. Then a hydrodynamic system $u_t^i = V_j^i u_x^j = (V^i)_{,j} u_x^j$ is compatible with A if and only if

$$\begin{cases} g_{im} V_j^m = g_{jm} V_i^m \\ c_{mkj} V_i^m + c_{mik} V_j^m + c_{mji} V_k^m = 0, \\ V_{i,j}^k = g^{ks} c_{smj} V_i^m + g^{ks} c_{smi} V_j^m \end{cases} \quad (1)$$

Each system is identified with a *congruence of lines* (Agafonov, Ferapontov (1996-2001)) in \mathbb{P}^{n+1} with coordinates $[y^1, \dots, y^{n+2}]$

$$y^i = u^i y^{n+1} + V^i y^{n+2}$$

The congruence is *linear*: there are n linear relations between $u^i, V^i, u^i V^j - u^j V^i$, hence the system is *linearly degenerate*, and *non diagonalizable*.

Digression: classification of operators and systems

(Results by Ferapontov, Pavlov, V.): The Hamiltonian operator

$$A^{ij} = \partial_x \left(g^{ij} \partial_x + c_k^{ij} u_x^k \right) \partial_x$$

is **completely determined** by its leading term g^{ij} . A multiparameter family of hydrodynamic systems is determined by a single operator A . Both objects are invariant up to reciprocal transformations of projective type:

$$\begin{aligned} d\tilde{x} &= (a_i u^i + a) dx + (a_i V^i + b) dt \\ d\tilde{t} &= (b_i u^i + c) dx + (b_i V^i + d) dt \end{aligned}$$

Classification: for $n = 3$ only two operators (and systems) exist, WDVV and a linear system. For $n = 4$ a classification is still possible; not possible for $n \geq 5$.

Results for N=3

$$\begin{cases} a_t = b_x, \\ b_t = c_x, \\ c_t = (b^2 - ac)_x \end{cases} \implies g_{ij} = \begin{pmatrix} -2b & a & 1 \\ a & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix} \quad \text{known}$$

$$\begin{cases} a_t = b_x, \\ b_t = c_x, \\ c_t = \left(\frac{-ac + b^2 + c^2 - 1}{b} \right)_x \end{cases} \implies g_{ij} = \begin{pmatrix} b^2 + 1 & -ab + bc & -b^2 \\ -ab + bc & a^2 - 2ac + c^2 + 1 & ab - bc \\ -b^2 & ab - bc & b^2 \end{pmatrix} \quad \text{new!}$$

The above system is indeed bi-Hamiltonian by the above operator and a first-order non-local operator (P. Lorenzoni, V. 2018).

$$\begin{aligned}
 & f_{2xt}^2(\eta_{11}\eta_{33} - \eta_{13}^2) + f_{2xt}f_{3t}(-\eta_{11}\eta_{22} + \eta_{12}^2) + f_{2xt}f_{x2t}(-\eta_{11}\eta_{23} + \eta_{12}\eta_{13}) \\
 & + f_{2xt}(-3\eta_{12}\eta_{23}\eta_{33} + \eta_{13}\eta_{22}\eta_{33} + 2\eta_{13}\eta_{23}^2) + f_{3t}f_{3x}(\eta_{11}\eta_{23} - \eta_{12}\eta_{13}) \\
 & + f_{3t}\eta_{22}(-\eta_{12}\eta_{23} + \eta_{13}\eta_{22}) + f_{3x}f_{x2t}(-\eta_{11}\eta_{33} + \eta_{13}^2) \\
 & + f_{3x}\eta_{33}(\eta_{12}\eta_{33} - \eta_{13}\eta_{23}) + f_{x2t}^2(\eta_{11}\eta_{22} - \eta_{12}^2) \\
 & + f_{x2t}(\eta_{12}\eta_{22}\eta_{33} + 2\eta_{12}\eta_{23}^2 - 3\eta_{13}\eta_{22}\eta_{23}) \\
 & - \eta_{22}^2\eta_{33}^2 + 2\eta_{22}\eta_{23}^2\eta_{33} - \eta_{23}^4 = 0 \quad (2)
 \end{aligned}$$

Theorem: The hydrodynamic system corresponding to the above WDVV equation yields a linear line congruence that is determined by a homogeneous third-order Hamiltonian operator. *Proof by computer.*

WDVV in the case $N = 4$

$$\eta = \begin{pmatrix} 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \end{pmatrix}$$

WDVV equations are an overdetermined nonlinear system:

$$\begin{aligned} -2f_{xyz} - f_{xyy}f_{xxy} + f_{yyy}f_{xxx} &= 0, \\ -f_{xzz} - f_{xyy}f_{xxz} + f_{yyz}f_{xxx} &= 0, \\ -2f_{xyz}f_{xxz} + f_{xzz}f_{xxy} + f_{yzz}f_{xxx} &= 0, \\ f_{zzz} - (f_{xyz})^2 + f_{xzz}f_{xyy} - f_{yyz}f_{xxz} + f_{yzz}f_{xxy} &= 0, \\ f_{yyy}f_{xzz} - 2f_{yyz}f_{xyz} + f_{yzz}f_{xyy} &= 0. \end{aligned}$$

6-components WDVV system

We introduce new field variables u^k :

$$u^1 = f_{xxx}, u^2 = f_{xxy}, u^3 = f_{xxz}, u^4 = f_{xyy}, u^5 = f_{xyz}, u^6 = f_{xzz}.$$

The compatibility conditions for this system can be written as a pair of *commuting* hydrodynamic type systems in conservative form:

$$\left\{ \begin{array}{l} u_y^1 = u_x^2, \\ u_y^2 = u_x^4, \\ u_y^3 = u_x^5, \\ u_y^4 = \left(\frac{2u^5 + u^2 u^4}{u^1} \right)_x \\ u_y^5 = \left(\frac{u^3 u^4 + u^6}{u^1} \right)_x \\ u_y^6 = \left(\frac{2u^3 u^5 - u^2 u^6}{u^1} \right)_x \end{array} \right. \quad \left\{ \begin{array}{l} u_z^1 = u_x^3, \\ u_z^2 = u_x^5, \\ u_z^3 = u_x^6, \\ u_z^4 = \left(\frac{u^3 u^4 + u^6}{u^1} \right)_x \\ u_z^5 = \left(\frac{2u^3 u^5 - u^2 u^6}{u^1} \right)_x \\ u_z^6 = \left(\frac{(u^5)^2 - u^4 u^6 + (u^3)^2 u^4 + u^3 u^6 - 2u^2 u^3 u^5 + (u^2)^2 u^6}{u^1} \right)_x \end{array} \right.$$

Theorem: (Pavlov, V. LMP 2015) The leading term of a third-order Hamiltonian operator for *both* the previous hydrodynamic-type systems:

$$g_{ik}(\mathbf{u}) = \begin{pmatrix} (u^4)^2 & -2u^5 & 2u^4 & -(u^1u^4 + u^3) & u^2 & 1 \\ -2u^5 & -2u^3 & u^2 & 0 & u^1 & 0 \\ 2u^4 & u^2 & 2 & -u^1 & 0 & 0 \\ -(u^1u^4 + u^3) & 0 & -u^1 & (u^1)^2 & 0 & 0 \\ u^2 & u^1 & 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

Both systems are bi-Hamiltonian with the above homogeneous third-order Hamiltonian operator and a (compatible) local first-order Hamiltonian operator.

Advanced computations

- ▶ In the case $N = 4$ also the choice $\eta_{\alpha\beta} = \delta_{\alpha\beta}$ yields a hydrodynamic system with a homogeneous third-order operator. It is, however, not so easy to prove that there exists a first-order operator.
- ▶ For generic N we obtain $N - 1$ hydrodynamic systems. Work is in progress in the case $N = 5$ to find a homogeneous third-order Hamiltonian operator.

Very advanced computations

The equation of associativity for F -manifolds, or *Oriented Associativity Equation* (Hertling–Manin '98):

$$\frac{\partial^2 c^i}{\partial a^j \partial a^m} \frac{\partial^2 c^m}{\partial a^k \partial a^n} = \frac{\partial^2 c^i}{\partial a^k \partial a^m} \frac{\partial^2 c^m}{\partial a^j \partial a^n} \quad (3)$$

also shows similar features: in the simplest case $N = 3$ it can be transformed into a hydrodynamic system which admits a third-order *non-local* Hamiltonian operator (Casati, Ferapontov, Pavlov, V. '18):

$$A^{ij} = \partial_x (g^{ij} \partial_x + c_k^{ij} u_x^k + w_k^i u_x^k \partial_x^{-1} w_h^j u_x^h) \partial_x,$$

with **extremely complicated** coordinate expression.

Open questions

| Integrable PDEs | Projective Geometry |
|----------------------------------|----------------------------|
| Third-order Hamiltonian operator | Quadratic Line Complex |
| Hydrodynamic system | Linear Line Congruence |
| First-order Hamiltonian operator | ??? |

Open questions: What is the Projective Geometric meaning of a first-order operator that is compatible with the third-order operator and with the hydrodynamic system? And what is the Projective Geometric meaning of the compatibility between operators and systems?

THANK YOU!