

Metric Invariants of Spherical Harmonics

Valentin Lychagin

IPU, RAS, Moscow & UiTo, Tromso, Norway

September 22, 2021

Spherical Harmonics

- ① Spherical harmonics of degree k are smooth solutions of the Euler

$$xu_x + yu_y + zu_z - ku = 0,$$

and the Laplace

$$u_{xx} + u_{yy} + u_{zz} = 0,$$

equations

- ① Spherical harmonics of degree k are smooth solutions of the Euler

$$xu_x + yu_y + zu_z - ku = 0,$$

and the Laplace

$$u_{xx} + u_{yy} + u_{zz} = 0,$$

equations

- ② Denote by

$$\mathcal{E}^{(i)} \subset J^i, i = 1, 2, \dots, k$$

the corresponding equations and their prolongations.

- 1 Spherical harmonics of degree k are smooth solutions of the Euler

$$xu_x + yu_y + zu_z - ku = 0,$$

and the Laplace

$$u_{xx} + u_{yy} + u_{zz} = 0,$$

equations

- 2 Denote by

$$\mathcal{E}^{(i)} \subset J^i, i = 1, 2, \dots, k$$

the corresponding equations and their prolongations.

- 3 Lie group $\mathbf{SO}(3)$ is obvious symmetry group of these equations and all $\mathcal{E}^{(i)}$ are affine algebraic manifolds equipped with the algebraic $\mathbf{SO}(3)$ -action.

- 1 Denote by \mathbb{H}_k the vector space of harmonic polynomials of degree k i.e. the solution space of the Euler-Laplace equations. It is a **SO** (3)-module.

- 1 Denote by \mathbb{H}_k the vector space of harmonic polynomials of degree k i.e. the solution space of the Euler-Laplace equations. It is a $\mathbf{SO}(3)$ -module.
- 2 We say that a rational $\mathbf{SO}(3)$ -invariant function on \mathbb{H}_k is an *algebraic metric invariant of spheric harmonics*, having degree k .

- 1 Denote by \mathbb{H}_k the vector space of harmonic polynomials of degree k i.e. the solution space of the Euler-Laplace equations. It is a $\mathbf{SO}(3)$ -module.
- 2 We say that a rational $\mathbf{SO}(3)$ -invariant function on \mathbb{H}_k is an *algebraic metric invariant of spheric harmonics*, having degree k .
- 3 The field of algebraic invariants we denote by \mathcal{F}_k^a .

- 1 Denote by \mathbb{H}_k the vector space of harmonic polynomials of degree k i.e. the solution space of the Euler-Laplace equations. It is a $\mathbf{SO}(3)$ -module.
- 2 We say that a rational $\mathbf{SO}(3)$ -invariant function on \mathbb{H}_k is an *algebraic metric invariant of spheric harmonics*, having degree k .
- 3 The field of algebraic invariants we denote by \mathcal{F}_k^a .
- 4 We say that a rational $\mathbf{SO}(3)$ -invariant function on affine manifold $\mathcal{E}^{(i)}$ is a *differential metric invariant of spheric harmonics*, having order $\leq i$.

- 1 Denote by \mathbb{H}_k the vector space of harmonic polynomials of degree k i.e. the solution space of the Euler-Laplace equations. It is a $\mathbf{SO}(3)$ -module.
- 2 We say that a rational $\mathbf{SO}(3)$ -invariant function on \mathbb{H}_k is an *algebraic metric invariant of spheric harmonics*, having degree k .
- 3 The field of algebraic invariants we denote by \mathcal{F}_k^a .
- 4 We say that a rational $\mathbf{SO}(3)$ -invariant function on affine manifold $\mathcal{E}^{(i)}$ is a *differential metric invariant of spheric harmonics*, having order $\leq i$.
- 5 The field of differentatial invariants we denote by \mathcal{F}_k^d

How structure on Weyl algebra

- ① The Weyl algebra \mathbb{A}_3 is the associative algebra of linear differential operators on \mathbb{R}^3 with polynomial coefficients.

How structure on Weyl algebra

- 1 The Weyl algebra \mathbb{A}_3 is the associative algebra of linear differential operators on \mathbb{R}^3 with polynomial coefficients.
- 2 The Lie algebra $\mathfrak{sl}(2) \subset \mathbb{A}_3$, generated by the following operators

$$X_+ = \frac{r^2}{2}, H = \delta + \frac{3}{2}, X_- = \frac{\Delta}{2},$$

where

$$r^2 = x^2 + y^2 + z^2, \delta = x\partial_x + y\partial_y + z\partial_z, \Delta = \partial_x^2 + \partial_y^2 + \partial_z^2,$$

and operators (X_+, H, X_-) form the Weyl basis in $\mathfrak{sl}(2)$:

$$[H, X_+] = 2X_+, [H, X_-] = -2X_-, [X_-, X_+] = H.$$

How structure on Weyl algebra

- 1 The Weyl algebra \mathbb{A}_3 is the associative algebra of linear differential operators on \mathbb{R}^3 with polynomial coefficients.
- 2 The Lie algebra $\mathfrak{sl}(2) \subset \mathbb{A}_3$, generated by the following operators

$$X_+ = \frac{r^2}{2}, H = \delta + \frac{3}{2}, X_- = \frac{\Delta}{2},$$

where

$$r^2 = x^2 + y^2 + z^2, \delta = x\partial_x + y\partial_y + z\partial_z, \Delta = \partial_x^2 + \partial_y^2 + \partial_z^2,$$

and operators (X_+, H, X_-) form the Weyl basis in $\mathfrak{sl}(2)$:

$$[H, X_+] = 2X_+, [H, X_-] = -2X_-, [X_-, X_+] = H.$$

- 3 The Lie algebra $\mathfrak{so}(3) \subset \mathbb{A}_3$ generated by the angular momentum operators

$$L_z = x\partial_y - y\partial_x, L_y = x\partial_z - z\partial_x, L_x = y\partial_z - z\partial_y.$$

How structure on Weyl algebra-2

- 1 These Lie algebras mutually commute and the *universal enveloping algebra* $U(\mathfrak{sl}(2)) \subset A_3$ is the subalgebra of *so(3)-invariant operators* in \mathbb{A}_3 .

How structure on Weyl algebra-2

- 1 These Lie algebras mutually commute and the *universal enveloping algebra* $U(\mathfrak{sl}(2)) \subset A_3$ is the subalgebra of *so(3)-invariant operators* in \mathbb{A}_3 .
- 2 *Casimir operator* in Lie algebra *so(3)* is the orbital angular momentum operator

$$M = L_x^2 + L_y^2 + L_z^2$$

and it coincides with *Casimir operator* in Lie algebra *sl(2)* :

$$M = r^2 \Delta - \delta^2 - \delta.$$

How structure on Weyl algebra-2

- 1 These Lie algebras mutually commute and the *universal enveloping algebra* $U(\mathfrak{sl}(2)) \subset A_3$ is the subalgebra of *so(3)-invariant operators* in \mathbb{A}_3 .
- 2 *Casimir operator* in Lie algebra *so(3)* is the orbital angular momentum operator

$$M = L_x^2 + L_y^2 + L_z^2$$

and it coincides with *Casimir operator* in Lie algebra *sl(2)* :

$$M = r^2 \Delta - \delta^2 - \delta.$$

- 3 Operator M is also the spherical Laplace operator.

Harmonic polynomials

- ① The following sequence

$$0 \rightarrow \mathbb{H}_k \rightarrow \mathbb{P}_k \xrightarrow{\Delta} \mathbb{P}_{k-2} \rightarrow 0$$

is exact, and $\dim \mathbb{H}_k = 2k + 1$.

Harmonic polynomials

- 1 The following sequence

$$0 \rightarrow \mathbb{H}_k \rightarrow \mathbb{P}_k \xrightarrow{\Delta} \mathbb{P}_{k-2} \rightarrow 0$$

is exact, and $\dim \mathbb{H}_k = 2k + 1$.

- 2 *Splitting* \mathbb{P}_k : for any homogeneous polynomial $p_k \in \mathbb{P}_k$ there are (and unique) spheric harmonics $h_{k-2i} \in \mathbb{H}_{k-2i}$, $0 \leq i \leq \lfloor \frac{k}{2} \rfloor$, such that

$$p = \sum_{i=0}^{\lfloor \frac{k}{2} \rfloor} r^{2i} h_{k-2i}.$$

Harmonic polynomials

- ① The following sequence

$$0 \rightarrow \mathbb{H}_k \rightarrow \mathbb{P}_k \xrightarrow{\Delta} \mathbb{P}_{k-2} \rightarrow 0$$

is exact, and $\dim \mathbb{H}_k = 2k + 1$.

- ② *Splitting* \mathbb{P}_k : for any homogeneous polynomial $p_k \in \mathbb{P}_k$ there are (and unique) spheric harmonics $h_{k-2i} \in \mathbb{H}_{k-2i}$, $0 \leq i \leq \lfloor \frac{k}{2} \rfloor$, such that

$$p = \sum_{i=0}^{\lfloor \frac{k}{2} \rfloor} r^{2i} h_{k-2i}.$$

- ③ We have

$$M(h_k) = -k(k+1)h_k,$$

for all $h_k \in \mathbb{H}_k$.

Harmonic polynomials

- 1 The following sequence

$$0 \rightarrow \mathbb{H}_k \rightarrow \mathbb{P}_k \xrightarrow{\Delta} \mathbb{P}_{k-2} \rightarrow 0$$

is exact, and $\dim \mathbb{H}_k = 2k + 1$.

- 2 *Splitting* \mathbb{P}_k : for any homogeneous polynomial $p_k \in \mathbb{P}_k$ there are (and unique) spheric harmonics $h_{k-2i} \in \mathbb{H}_{k-2i}$, $0 \leq i \leq \lfloor \frac{k}{2} \rfloor$, such that

$$p = \sum_{i=0}^{\lfloor \frac{k}{2} \rfloor} r^{2i} h_{k-2i}.$$

- 3 We have

$$M(h_k) = -k(k+1)h_k,$$

for all $h_k \in \mathbb{H}_k$.

- 4 The restriction of spheric harmonics on the unit sphere $\mathbf{S}^2 \subset \mathbb{R}^3$ are eigenfunctions of the spherical laplacian Δ_S with eigenvalues $-k(k+1)$ and any continuous function on \mathbf{S}^2 could be approximated (with any accuracy) by linear combination of spherical harmonics.

Harmonic projections

- ① *Harmonic projections* $\eta_{k,2i} : \mathbb{P}_k \rightarrow \mathbb{H}_{k-2i}$ are the following

$$\eta_{k,2i} = r^{-2i} Q_{k,2i}(M),$$

where

$$Q_{k,2i}(\lambda) = \prod_{j \neq i}^{\lfloor \frac{k}{2} \rfloor} \frac{\lambda - \lambda_j}{\lambda_i - \lambda_j}, \quad \lambda_i = -(k - 2i)(k - 2i + 1).$$

Harmonic projections

- ① *Harmonic projections* $\eta_{k,2i} : \mathbb{P}_k \rightarrow \mathbb{H}_{k-2i}$ are the following

$$\eta_{k,2i} = r^{-2i} Q_{k,2i}(M),$$

where

$$Q_{k,2i}(\lambda) = \prod_{j \neq i}^{\lfloor \frac{k}{2} \rfloor} \frac{\lambda - \lambda_j}{\lambda_i - \lambda_j}, \quad \lambda_i = -(k - 2i)(k - 2i + 1).$$

- ② The following sequence

$$\mathbf{0} \rightarrow \mathbb{P}_{k-2} \xrightarrow{r^2} \mathbb{P}_k \xrightarrow{\eta_{k,0}} \mathbb{H}_k \rightarrow \mathbf{0}$$

is exact.

Harmonic projections

- ① Harmonic projections $\eta_{k,2i} : \mathbb{P}_k \rightarrow \mathbb{H}_{k-2i}$ are the following

$$\eta_{k,2i} = r^{-2i} Q_{k,2i}(M),$$

where

$$Q_{k,2i}(\lambda) = \prod_{j \neq i}^{\lfloor \frac{k}{2} \rfloor} \frac{\lambda - \lambda_j}{\lambda_i - \lambda_j}, \quad \lambda_i = -(k - 2i)(k - 2i + 1).$$

- ② The following sequence

$$\mathbf{0} \rightarrow \mathbb{P}_{k-2} \xrightarrow{r^2} \mathbb{P}_k \xrightarrow{\eta_{k,0}} \mathbb{H}_k \rightarrow \mathbf{0}$$

is exact.

- ③ Define product of spheric harmonics $h_k \in \mathbb{H}_k, h_l \in \mathbb{H}_l$ as follows

$$h_k * h_l = \eta_{k+l,0}(h_k h_l) \in \mathbb{H}_{k+l}.$$

Harmonic projections

- ① Harmonic projections $\eta_{k,2i} : \mathbb{P}_k \rightarrow \mathbb{H}_{k-2i}$ are the following

$$\eta_{k,2i} = r^{-2i} Q_{k,2i}(M),$$

where

$$Q_{k,2i}(\lambda) = \prod_{j \neq i}^{\lfloor \frac{k}{2} \rfloor} \frac{\lambda - \lambda_j}{\lambda_i - \lambda_j}, \quad \lambda_i = -(k - 2i)(k - 2i + 1).$$

- ② The following sequence

$$\mathbf{0} \rightarrow \mathbb{P}_{k-2} \xrightarrow{r^2} \mathbb{P}_k \xrightarrow{\eta_{k,0}} \mathbb{H}_k \rightarrow \mathbf{0}$$

is exact.

- ③ Define product of spheric harmonics $h_k \in \mathbb{H}_k, h_l \in \mathbb{H}_l$ as follows

$$h_k * h_l = \eta_{k+l,0}(h_k h_l) \in \mathbb{H}_{k+l}.$$

- ④ Here

$$\eta_{k+l,0} = \prod_{j=1}^{\lfloor \frac{k+l}{2} \rfloor} \frac{M + (k+l-2j)(k+l-2j+1)}{2j(2j-2k-2l-1)}$$

1 Spherical harmonics

$$\mathbb{H}_* = \bigoplus_{k \geq 0} \mathbb{H}_k$$

form a graded commutative algebra with respect to the product $*$.

1 Spherical harmonics

$$\mathbb{H}_* = \bigoplus_{k \geq 0} \mathbb{H}_k$$

form a graded commutative algebra with respect to the product $*$.

2 Algebra $(\mathbb{H}_*, *)$ generated by linear functions (x, y, z) satisfying the relation

$$x * x + y * y + z * z = 0.$$

Algebra of spherical harmonics

1 Spherical harmonics

$$\mathbb{H}_* = \bigoplus_{k \geq 0} \mathbb{H}_k$$

form a graded commutative algebra with respect to the product $*$.

2 Algebra $(\mathbb{H}_*, *)$ generated by linear functions (x, y, z) satisfying the relation

$$x * x + y * y + z * z = 0.$$

3 The complification $\mathbb{H}_* = \mathbb{H} \otimes \mathbb{C}$ is the algebra of regular functions on the null cone $\{x^2 + y^2 + z^2 = 0\}$ in \mathbb{C}^3 .

Algebra of spherical harmonics

1 Spherical harmonics

$$\mathbb{H}_* = \bigoplus_{k \geq 0} \mathbb{H}_k$$

form a graded commutative algebra with respect to the product $*$.

2 Algebra $(\mathbb{H}_*, *)$ generated by linear functions (x, y, z) satisfying the relation

$$x * x + y * y + z * z = 0.$$

3 The complexification $\mathbb{H}_* = \mathbb{H} \otimes \mathbb{C}$ is the algebra of regular functions on the null cone $\{x^2 + y^2 + z^2 = 0\}$ in \mathbb{C}^3 .

4 Example.

$$x * x = xx - \frac{r^2}{3}, x * y = xy.$$

- 1 The space \mathbb{H}_k of spherical harmonics is a vector space of dimension $2k + 1$. The Lie group $\mathbf{SO}(3)$ acts in algebraic way on \mathbb{H}_k , and in \mathbb{H}_k are realized all irreducible representations of $\mathbf{SO}(3)$.

Algebraic invariants

- 1 The space \mathbb{H}_k of spherical harmonics is a vector space of dimension $2k + 1$. The Lie group $\mathbf{SO}(3)$ acts in algebraic way on \mathbb{H}_k , and in \mathbb{H}_k are realized all irreducible representations of $\mathbf{SO}(3)$.
- 2 Due to Hilbert theorem polynomial invariants of this action (i.e. polynomial invariants of spherical harmonics) form a finite generated commutative algebra.

Algebraic invariants

- 1 The space \mathbb{H}_k of spherical harmonics is a vector space of dimension $2k + 1$. The Lie group $\mathbf{SO}(3)$ acts in algebraic way on \mathbb{H}_k , and in \mathbb{H}_k are realized all irreducible representations of $\mathbf{SO}(3)$.
- 2 Due to Hilbert theorem polynomial invariants of this action (i.e. polynomial invariants of spherical harmonics) form a finite generated commutative algebra.
- 3 Due to Rosenlicht theorem rational invariants of this action (i.e. rational invariants of spherical harmonics) form a field of transcendence degree equals the codimension of regular orbit.

- 1 The space \mathbb{H}_k of spherical harmonics is a vector space of dimension $2k + 1$. The Lie group $\mathbf{SO}(3)$ acts in algebraic way on \mathbb{H}_k , and in \mathbb{H}_k are realized all irreducible representations of $\mathbf{SO}(3)$.
- 2 Due to Hilbert theorem polynomial invariants of this action (i.e. polynomial invariants of spherical harmonics) form a finite generated commutative algebra.
- 3 Due to Rosenlicht theorem rational invariants of this action (i.e. rational invariants of spherical harmonics) form a field of transcendence degree equals the codimension of regular orbit.
- 4 Regular orbit has codimension $(2k - 2)$, when $k \geq 2$, and codimension 1, when $k = 1$. Therefore, in order to define a regular orbit we need $2k - 2$ algebraically independent rational invariants, for $k > 2$, and only one invariant, for $k = 1$.

- Equations $\mathcal{E}^{(i)}$ are affine manifolds of dimension $2i + 4$, if $2 \leq i < k$. The regular $\mathbf{SO}(3)$ -orbits (that correspond to smooth points of quotient $\mathcal{E}^{(i)}/\mathbf{SO}(3)$) Thus, due to Hilbert theorem, the quotients are affine manifolds of dimension $2i + 1$. Rational differential invariants of order $\leq i$ are rational functions on $\mathcal{E}^{(i)}/\mathbf{SO}(3)$ and therefore the transcendence degree of field \mathcal{F}_i^d equals to $2i + 1$.

Differential invariants

- 1 Equations $\mathcal{E}^{(i)}$ are affine manifolds of dimension $2i + 4$, if $2 \leq i < k$. The regular $\mathbf{SO}(3)$ -orbits (that correspond to smooth points of quotient $\mathcal{E}^{(i)}/\mathbf{SO}(3)$) Thus, due to Hilbert theorem, the quotients are affine manifolds of dimension $2i + 1$. Rational differential invariants of order $\leq i$ are rational functions on $\mathcal{E}^{(i)}/\mathbf{SO}(3)$ and therefore the transcendence degree of field \mathcal{F}_i^d equals to $2i + 1$.
- 2 As we have seen, the transcendence degree of field \mathcal{F}_k^a equals $2(k - 1)$.

- 1 Equations $\mathcal{E}^{(i)}$ are affine manifolds of dimension $2i + 4$, if $2 \leq i < k$. The regular $\mathbf{SO}(3)$ -orbits (that correspond to smooth points of quotient $\mathcal{E}^{(i)}/\mathbf{SO}(3)$) Thus, due to Hilbert theorem, the quotients are affine manifolds of dimension $2i + 1$. Rational differential invariants of order $\leq i$ are rational functions on $\mathcal{E}^{(i)}/\mathbf{SO}(3)$ and therefore the transcendence degree of field \mathcal{F}_i^d equals to $2i + 1$.
- 2 As we have seen, the transcendence degree of field \mathcal{F}_k^a equals $2(k - 1)$.
- 3 Take a regular harmonic $h \in \mathbb{H}_k$. Then it is easy to check that the $\mathbf{SO}(3)$ -orbit of the 2-jet $j_2(h)$ into $\mathcal{E}^{(2)}$ is a 6-dimensional submanifold into 8-dimensional manifold $\mathcal{E}^{(2)}$ and therefore *we need 2 differential invariants of order 2 to describe the orbit* (compare with $2(k - 1)$ algebraic invariants).

linvariant coframe

- ① Total differentials of the obvious invariants $J_{-1} = \frac{r^2}{2}$ and $J_0 = u$ give us two **SO**(3)-invariant horizontal 1-forms:

$$\omega_1 = xdx + ydy + zdz,$$

$$\omega_2 = u_x dx + u_y dy + u_z dz.$$

linvariant coframe

- ① Total differentials of the obvious invariants $J_{-1} = \frac{r^2}{2}$ and $J_0 = u$ give us two **SO**(3)-invariant horizontal 1-forms:

$$\omega_1 = xdx + ydy + zdz,$$

$$\omega_2 = u_x dx + u_y dy + u_z dz.$$

- ② Their cross product gives us

$$\omega_3 = (yu_z - zu_y) dx + (zu_x - xu_z) dy + (xu_y - yu_x) dz.$$

linvariant coframe

- ① Total differentials of the obvious invariants $J_{-1} = \frac{r^2}{2}$ and $J_0 = u$ give us two **SO**(3)-invariant horizontal 1-forms:

$$\omega_1 = xdx + ydy + zdz,$$

$$\omega_2 = u_x dx + u_y dy + u_z dz.$$

- ② Their cross product gives us

$$\omega_3 = (yu_z - zu_y) dx + (zu_x - xu_z) dy + (xu_y - yu_x) dz.$$

- ③ Then coframe $(\omega_1, \omega_2, \omega_3)$ is **SO**(3)-invariant.

Invariant frame

1

$$D_1 = x \frac{d}{dx} + y \frac{d}{dy} + z \frac{d}{dz},$$

$$D_2 = u_x \frac{d}{dx} + u_y \frac{d}{dy} + u_z \frac{d}{dz},$$

$$D_3 = (yu_z - zu_y) \frac{d}{dx} + (zu_x - xu_z) \frac{d}{dy} + (xu_y - yu_x) \frac{d}{dz}.$$

First invariants

$$J_{-1} = \frac{r^2}{2}, J_0 = u,$$

$$J_1 = D_2(J_0) = u_x^2 + u_y^2 + u_z^2,$$

$$J_{21} = \frac{D_2(J_1)}{2} = u_x^2 u_{xx} + u_y^2 u_{yy} + u_z^2 u_{zz} + 2(u_x u_y u_{xy} + u_x u_z u_{xz} + u_y u_z u_{yz}).$$

Invariant symmetric forms and operators

1 Symmetric differential i -forms

$$\Theta_i = \sum_{i_1+i_2+i_3=i} u_{i_1, i_2, i_3} \frac{dx^{i_1}}{i_1!} \cdot \frac{dy^{i_2}}{i_2!} \cdot \frac{dz^{i_3}}{i_3!}$$

are invariants with respect to Lie group of affine transformations in \mathbb{R}^3 .

Invariant symmetric forms and operators

1 Symmetric differential i -forms

$$\Theta_i = \sum_{i_1+i_2+i_3=i} u_{i_1,i_2,i_3} \frac{dx^{i_1}}{i_1!} \cdot \frac{dy^{i_2}}{i_2!} \cdot \frac{dz^{i_3}}{i_3!}$$

are invariants with respect to Lie group of affine transformations in \mathbb{R}^3 .

2 Differential operators

$$\widehat{\Theta}_i = \sum_{i_1+i_2+i_3=i} \frac{u_{i_1,i_2,i_3}}{i_1!i_2!i_3!} \frac{d^k}{dx^{i_1} dy^{i_2} dz^{i_3}}$$

are **SO**(3)-invariant.

Invariants

Let

$$dx = t_{11}\omega_1 + t_{12}\omega_2 + t_{13}\omega_3,$$

$$dy = t_{21}\omega_1 + t_{22}\omega_2 + t_{23}\omega_3,$$

$$dz = t_{31}\omega_1 + t_{32}\omega_2 + t_{33}\omega_3,$$

where t_{ij} are rational functions on $J^1(\mathbb{R}^3)$, and let

$$\Theta_i = \sum_{i_1+i_2+i_3=i} T_{i_1, i_2, i_3} \frac{\omega_1^{i_1}}{i_1!} \cdot \frac{\omega_2^{i_2}}{i_2!} \cdot \frac{\omega_3^{i_3}}{i_3!}.$$

Theorem

Functions T_{i_1, i_2, i_3} are rational differential $\mathbf{SO}(3)$ -invariants of order $i = i_1 + i_2 + i_3$ and any rational differential $\mathbf{SO}(3)$ -invariants of order i is a rational function of them.

Example

Remark that invariants

$$G_i = \widehat{\Theta}_i(u) = \sum_{i_1+i_2+i_3=i} \frac{u_{i_1, i_2, i_3}^2}{i_1! i_2! i_3!}$$

are squares of lengths of symmetric forms Θ_i .

Thus,

$$\widehat{\Theta}_1 = u_x \frac{d}{dx} + u_y \frac{d}{dy} + u_z \frac{d}{dz},$$

$$\widehat{\Theta}_2 = \frac{1}{2} \left(u_{xx} \frac{d^2}{dx^2} + u_{yy} \frac{d^2}{dy^2} + u_{zz} \frac{d^2}{dz^2} \right) + u_{xy} \frac{d^2}{dxdy} + u_{xz} \frac{d^2}{dxdz} + u_{yz} \frac{d^2}{dydz}$$

and

$$\widehat{\Theta}_1(u) = u_x^2 + u_y^2 + u_z^2,$$

$$\widehat{\Theta}_2(u) = J_{22} = \frac{u_{xx}^2 + u_{yy}^2 + u_{zz}^2}{2} + u_{xy}^2 + u_{xz}^2 + u_{yz}^2.$$

Theorem

The field of rational differential $\mathbf{SO}(3)$ -invariants of spherical harmonics is generated by invariants $\left(J_{-1} = \frac{r^2}{2}, J_0 = u, J_{22}\right)$ and derivation $\nabla = \widehat{\mathcal{H}}_1$.

① $SO(3)$ –Invariants $\iff SO(3)$ –invariant differential operators:

$$\begin{aligned} \phi \in C^\infty(\mathbf{J}^k(\mathbb{R}^3)) &\iff \Delta_\phi : C^\infty(\mathbb{R}^3) \rightarrow C^\infty(\mathbb{R}^3), \\ \Delta_\phi(f) &= j_k(f)^*(\phi). \end{aligned}$$

- ① **SO**(3) –Invariants \iff **SO**(3) –invariant differential operators:

$$\phi \in C^\infty(\mathbf{J}^k(\mathbb{R}^3)) \iff \Delta_\phi : C^\infty(\mathbb{R}^3) \rightarrow C^\infty(\mathbb{R}^3),$$
$$\Delta_\phi(f) = j_k(f)^*(\phi).$$

- ② Monoid structure on **SO**(3) –invariants defines by the composition of invariant operators, and $\text{id} = u$.

- ① $\mathbf{SO}(3)$ –Invariants $\iff SO(3)$ –invariant differential operators:

$$\phi \in C^\infty(\mathbf{J}^k(\mathbb{R}^3)) \iff \Delta_\phi : C^\infty(\mathbb{R}^3) \rightarrow C^\infty(\mathbb{R}^3),$$
$$\Delta_\phi(f) = j_k(f)^*(\phi).$$

- ② Monoid structure on $\mathbf{SO}(3)$ –invariants defines by the composition of invariant operators, and $\text{id} = u$.
- ③ Thus, the field \mathcal{F}_k^k is the monoid.

① Let

$$W = x\partial_x + y\partial_y + z\partial_z + ku\partial_u,$$

and let W^* be its ∞ -prolongation.

- 1 Let

$$W = x\partial_x + y\partial_y + z\partial_z + ku\partial_u,$$

and let W^* be its ∞ -prolongation.

- 2 We say that a polynomial differential invariant I has weight $w(I)$ if

$$W^*(I) = w(I)I.$$

- 1 Let

$$W = x\partial_x + y\partial_y + z\partial_z + ku\partial_u,$$

and let W^* be its ∞ -prolongation.

- 2 We say that a polynomial differential invariant I has weight $w(I)$ if

$$W^*(I) = w(I)I.$$

- 3 In other words, if h is a homogeneous polynomial of degree k then $I(h)$ has degree $w(I)$.

- 1 Algebraic invariants on \mathbb{H}_k are differential invariants of order k .

Differential or Algebraic invariants

- 1 Algebraic invariants on \mathbb{H}_k are differential invariants of order k .
- 2 Let I be a polynomial differential invariant of weight w , and $h \in \mathbb{H}_k$.

- 1 Algebraic invariants on \mathbb{H}_k are differential invariants of order k .
- 2 Let I be a polynomial differential invariant of weight w , and $h \in \mathbb{H}_k$.
- 3 Then $I(h) \in \mathbb{P}_w$, $(\eta_{w,2I} \circ \Delta_I)(h) \in \mathbb{H}_{w-2I}$ and its length $(\Delta_{G_{w-2I}} \circ \eta_{w,2I} \circ \Delta_I)(h)$ is a scalar, i.e invariant $G_{w-2I} \circ \eta_{w,2I} \circ I$ is an algebraic invariant.