

Reciprocal transformations and deformations of integrable hierarchies

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***R*-matrix formalism**

A systematic way to construct the operators A_k for a given L so that we obtain self-consistent compatibility conditions $L_{t_k} = [A_k, L]$ is provided by the *R-matrix formalism*.

Given a Lie algebra \mathfrak{G} with the commutator $[\cdot, \cdot]$, a linear map $R : \mathfrak{G} \rightarrow \mathfrak{G}$ such that the bracket

$$[a, b]_R := [Ra, b] + [a, Rb] \quad (1)$$

is another Lie bracket on \mathfrak{G} is called the classical *R*-matrix.

***R*-matrix formalism II**

A sufficient condition for R to be a classical R -matrix is to satisfy the following so-called Yang-Baxter equation, YB(α),

$$[Ra, Rb] - R[a, b]_R + \alpha [a, b] = 0, \quad (2)$$

where α is a number from the ground field \mathbb{K} (\mathbb{R} or \mathbb{C}). There are only two substantially different cases, namely $\alpha \neq 0$ and $\alpha = 0$.

Algebra of pseudo-differential operators

$$\mathfrak{A} = \left\{ L = \sum_{i=-\infty}^N h_i D^i \right\}, \quad (3)$$

where D is (informally) the total x -derivative, where x is the independent variable.

The noncommutative multiplication in \mathfrak{A} is defined using the (generalized) Leibniz rule

$$D^k \circ f = \sum_{i=0}^{\infty} \frac{k(k-1)\cdots(k-i+1)}{i!} D^i(f) D^{k-i} \quad (4)$$

$$[A, B] := A \circ B - B \circ A \quad A, B \in \mathfrak{A}.$$

Algebra of pseudo-differential operators II

Consider the following decomposition of \mathfrak{A} :

$$\mathfrak{A} = \mathfrak{A}_{\geq k} \oplus \mathfrak{A}_{< k} := \left\{ \sum_{i \geq k} u_i D^i \right\} \oplus \left\{ \sum_{i < k} a_i D^i \right\}. \quad (5)$$

Then, $\mathfrak{A}_{\geq k}$ and $\mathfrak{A}_{< k}$ are Lie subalgebras of \mathfrak{A} only for $k = 0, 1, 2$, and we have the classical R -matrices

$$R_k = \frac{1}{2}(P_{\geq k} - P_{< k}) = P_{\geq k} - \frac{1}{2} = \frac{1}{2} - P_{< k}, \quad (6)$$

where $P_{\geq k}$ and $P_{< k}$ are projections onto $\mathfrak{A}_{\geq k}$ and $\mathfrak{A}_{< k}$, respectively.

Fractional powers

Consider an element L from \mathfrak{A} of the form

$$L = u_N D^N + u_{N-1} D^{N-1} + u_{N-2} D^{N-2} + \dots, \quad (7)$$

where $N > 0$. Then its N -th root

$$L^{\frac{1}{N}} = a_1 D + a_0 + a_{-1} D^{-1} + a_{-2} D^{-2} + \dots,$$

where the coefficients a_i are **local** (depend on u_i and a finite number of their x -derivatives) and can be constructed recursively from the equality

$$\left(L^{\frac{1}{N}}\right)^N = \overbrace{L^{\frac{1}{N}} \circ \dots \circ L^{\frac{1}{N}}}^{N \text{ times}} = L.$$

Fractional powers and Lax hierarchies on \mathfrak{A}

The fractional powers $L^{\frac{n}{N}} = \overbrace{L^{\frac{1}{N}} \circ \dots \circ L^{\frac{1}{N}}}^{n \text{ times}}$ generate the following Lax hierarchies related to classical R -matrices (6):

$$L_{t_n} = \left[R_k(L^{\frac{n}{N}}), L \right] = \left[\left(L^{\frac{n}{N}} \right)_{\geq k}, L \right] = - \left[\left(L^{\frac{n}{N}} \right)_{< k}, L \right],$$

$$k = 0, 1, 2, \quad n = 1, 2, \dots, \quad n \not\equiv 0 \pmod{N} \quad (8)$$

Key result: commutativity of the flows:

$$(L_{t_n})_{t_m} = (L_{t_m})_{t_n}$$

follows from the YBE for R_k .

More on Lax hierarchies on \mathfrak{A}

The Lax hierarchies involving *finitely many* dynamical variables u_i are obtained within this approach by choosing a special form of L .

$$k = 0, 1, 2 : L = \sum_{j=0}^N u_j D^j, \quad (9)$$

$$k = 1, 2 : L = \sum_{j=0}^N u_j D^j + D^{-1} \circ u_{-1}, \quad (10)$$

$$k = 2 : L = \sum_{j=0}^N u_j D^j + D^{-1} \circ u_{-1} + D^{-2} \circ u_{-2}, \quad (11)$$

It can be shown that without loss of generality we can set $u_N = 1$ in (9) and (10).

Is there more than that?

Błaszak and Szablikowski (*J. Math. Phys.* 2006) found the following deformations of the above R -matrices R_k :

$$R'_k = P_{\geq k} - \frac{1}{2} + \epsilon P_{k-1}(\cdot) D^k = \frac{1}{2} - P_{< k} + \epsilon P_{k-1}(\cdot) D^k, \text{ i.e.,}$$

$$\begin{aligned} R'_k(L) &= P_{\geq k}(L) - \frac{1}{2}L + \epsilon P_{k-1}(L) D^k \\ &= \frac{1}{2}L - P_{< k}(L) + \epsilon P_{k-1}(L) D^k, \quad L \in \mathfrak{A}. \end{aligned}$$

Here $P_k = P_{\geq k} - P_{\geq k-1}$ and ϵ is an arbitrary constant.

Deformed Lax hierarchies

$$\begin{aligned}
 L_{t_n} &= \left[R'_k \left(L^{\frac{n}{N}} \right), L \right] = \left[\left(L^{\frac{n}{N}} \right)_{\geq k} + \epsilon P_{k-1} \left(L^{\frac{n}{N}} \right) D^k, L \right] \\
 &= - \left[\left(L^{\frac{n}{N}} \right)_{< k} - \epsilon P_{k-1} \left(L^{\frac{n}{N}} \right) D^k, L \right], \\
 k &= 0, 1, 2, \quad n = 1, 2, \dots, \quad n \not\equiv 0 \pmod{N}
 \end{aligned} \tag{12}$$

However, for $k = 0$ and $k = 2$ Błaszak & Szablikowski have shown that these hierarchies can be reduced to known ones, i.e., the deformations in question are trivial.

Problem: is this true also for $k = 1$?

Trivializing transformation for $k = 1$

Consider the *reciprocal transformation* from x and t_i , $i = 1, 2, \dots, i \not\equiv 0 \pmod{N}$ to new independent variables z and τ_i , $i = 1, 2, \dots$, where $\tau_i = t_i$, $i = 1, 2, \dots, i \not\equiv 0 \pmod{N}$ and z is defined by the formula

$$dz = (u_N)^{-1/N} dx + \epsilon \sum_{q=1, q \not\equiv 0 \pmod{N}}^{\infty} (u_N)^{-1/N} P_0(L^{q/N}) dt_q. \quad (13)$$

Trivializing transformation for $k = 1$: continued

Introduce new dependent variables v_i related to u_i by means of the formulas

$$\begin{aligned}\tilde{L} &= L|_{D=(u_N)^{-1/N}\tilde{D}} = \sum_{i=-\infty}^N u_i ((u_N)^{-1/N} \tilde{D})^i \\ &\equiv \tilde{D}^N + \sum_{i=-\infty}^{N-1} v_i \tilde{D}^i\end{aligned}\tag{14}$$

(informally, \tilde{D} is the total z -derivative) and

$$v_N = (u_N)^{1/N}.\tag{15}$$

Main result

Theorem 1 *The above transformation sends the hierarchy*

$$L_{t_n} = \left[R'_1(L^{\frac{n}{N}}, L) \right] = \left[\left(L^{\frac{n}{N}} \right)_{\geq 1} + \epsilon P_0(L^{\frac{n}{N}}) D, L \right]$$

$$k = 0, 1, 2, \quad n = 1, 2, \dots, \quad n \not\equiv 0 \pmod{N}$$

into the undeformed $k = 1$ -hierarchy for v_i ,

$$\tilde{L}_{\tau_q} = [P_{\geq 1}(\tilde{L}^q), \tilde{L}], \quad q = 1, 2, \dots, q \not\equiv 0 \pmod{N},$$

for $\tilde{L} = \tilde{D}^N + \sum_{i=-\infty}^{N-1} v_i \tilde{D}^i$ along with a hierarchy of equations for v_N :

$$(v_N)_{\tau_q} = -\epsilon \tilde{D}(P_0(\tilde{L}^{q/N})), \quad q = 1, 2, \dots, q \not\equiv 0 \pmod{N}.$$

Example: Extended Broer–Kaup hierarchy

Consider the extended Broer–Kaup system with

$$L = uD + v + D^{-1} \circ w. \text{ From}$$

$$L_{t_i} = [P_{\geq 1}(L^i) + \epsilon P_0(L^i)D, L]$$

with $i = 1, 2$ we obtain

$$\begin{pmatrix} u \\ v \\ w \end{pmatrix}_{t_1} = \begin{pmatrix} \epsilon u_x v - \epsilon u v_x \\ u v_x + \epsilon v v_x \\ u_x w + u w_x + \epsilon v_x w + \epsilon v w_x \end{pmatrix}$$

Extended Broer–Kaup system

$$\begin{aligned}u_{t_2} &= \epsilon u_x v^2 - 2\epsilon u v v_x - 2\epsilon u^2 w_x - \epsilon u^2 v_{xx} \\v_{t_2} &= 2u u_x w + 2u v v_x + 2u^2 w_x + u u_x v_x + u^2 v_{xx} \\&\quad + \epsilon v^2 v_x + 2\epsilon u v_x w + \epsilon u v_x^2 \\w_{t_2} &= 2u_x v w + 2u v_x w + 2u v w_x - u_x^2 w - 3u u_x w_x \\&\quad - u u_{xx} w - u^2 w_{xx} + 2\epsilon u_x w^2 \\&\quad + 2\epsilon v v_x w + \epsilon u_x v_x w + \epsilon v^2 w_x + 4\epsilon u w w_x \\&\quad + \epsilon u v_x w_x + \epsilon u v_{xx} w.\end{aligned}$$

Upon setting $\epsilon = 0$ and $u = 1$ we recover from the second (t_2) flow the standard Kaup–Broer system.

Transformation to the Broer–Kaup hierarchy

By Theorem 1, pass from x and $t_{1,2}$ to z and $\tau_{1,2}$ defined by the formulas $\tau_1 = t_1$, $\tau_2 = t_2$, and

$$dz = (1/u)dx + \epsilon(v/u)dt_1 + (\epsilon/u)(2wu + uv_x + v^2)dt_2$$

(we ignore here the times t_i with $i \neq 1, 2$).

We have $\partial_x = (1/u)\partial_z$, so L goes into

$$\tilde{L} = \left(uD + v + D^{-1} \circ w \right)_{D=(1/u)\tilde{D}} = \tilde{D} + v + \tilde{D}^{-1} \circ r,$$

where $r = wu$.

Transformation to the Broer–Kaup hierarchy II

We readily see that the transformed hierarchy

$$\tilde{L}_{\tau_q} = [P_{\geq 1}(\tilde{L}^q), \tilde{L}], \quad q = 1, 2, \dots,$$

is nothing but the *standard* Broer–Kaup hierarchy for v, r with the independent variables z and τ_i , and we have separated equations for u which are linear on the background of this hierarchy:

$$u_{\tau_q} = -\epsilon \tilde{D}(P_0(\tilde{L}^q/N)), \quad q = 1, 2, \dots$$

Dispersionless limit

Theorem 1 remains valid if we replace the algebra \mathfrak{A} by its dispersionless (or *quasiclassical*) limit:

$D \rightarrow p$, $[,] \rightarrow \{, \}$, $\{, \}$ is the Poisson bracket:

$$\{f, g\} = D(f) \frac{\partial g}{\partial p} - D(g) \frac{\partial f}{\partial p}.$$

Now $L = \sum_{i=-\infty}^N u_i p^i$ is a function rather than an operator, the reciprocal transformation is the same: from x and t_i we pass to z and τ_i , where

$\tau_i = t_i$, $i = 1, 2, \dots$, where z is defined by

$$dz = (u_N)^{-1/N} dx + \epsilon \sum_{q=1, q \not\equiv 0 \pmod{N}}^{\infty} (u_N)^{-1/N} P_0(L^{q/N}) dt_q.$$

Dispersionless limit II

New dependent variables v_i are now related to u_i by means of the formulas

$$\begin{aligned}\tilde{L} &= L|_{p=(u_N)^{-1/N}\tilde{p}} = \sum_{i=-\infty}^N u_i ((u_N)^{-1/N}\tilde{p})^i \\ &\equiv \tilde{p}^N + \sum_{i=-\infty}^{N-1} v_i \tilde{p}^i,\end{aligned}$$

i.e.,

$$v_i = u_i (u_N)^{-i/N}, \quad (16)$$

and

$$v_N = (u_N)^{1/N}. \quad (17)$$

Main result for the dispersionless case

Theorem 2 *The above transformation sends*

$$L_{t_n} = \{R'_1(L^{\frac{n}{N}}), L\} = \left\{ \left(L^{\frac{n}{N}} \right)_{\geq 1} + \epsilon P_0(L^{\frac{n}{N}})p, L \right\}$$

$$k = 0, 1, 2, \quad n = 1, 2, \dots, \quad n \not\equiv 0 \pmod{N}$$

into the undeformed $k = 1$ -hierarchy for v_i ,

$$\tilde{L}_{\tau_q} = \{P_{\geq 1}(\tilde{L}^q), \tilde{L}\}', \quad q = 1, 2, \dots, q \not\equiv 0 \pmod{N},$$

for $\tilde{L} = \tilde{p}^N + \sum_{i=-\infty}^{N-1} v_i \tilde{p}^i$, where

$$\{f, g\}' = \tilde{D}(f) \frac{\partial g}{\partial \tilde{p}} - \tilde{D}(g) \frac{\partial f}{\partial \tilde{p}},$$

along with a hierarchy for v_N :

$$(v_N)_{\tau_q} = -\epsilon \tilde{D}(P_0(\tilde{L}^{q/N})), \quad q = 1, 2, \dots, q \not\equiv 0 \pmod{N}.$$

Conclusions

All integrable hierarchies constructed using the deformed R -matrices on \mathfrak{A} given for $k = 0, 1, 2$ by

$$\begin{aligned} R'_k(L) &= P_{\geq k}(L) - \frac{1}{2}L + \epsilon P_{k-1}(L)D^k \\ &= \frac{1}{2}L - P_{< k}(L) + \epsilon P_{k-1}(L)D^k, \quad L \in \mathfrak{A}, \end{aligned}$$

can be reduced to linear extensions of the undeformed ($\epsilon = 0$) hierarchies using suitable changes of dependent and independent variables. The same holds for their dispersionless counterparts.