Metrics admitting projective symmetries

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This and the encouragement to think "simply" are the most important teachings I got from him.

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The class of connections sharing the same geodesics of a connection  $\Gamma$  is called *projective connection* (associated to  $\Gamma$ ).

### Problems of Sophus Lie

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Untersuchungen über geodätische Curven

(Math. Ann. 20, 1882)

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II. Man soll die Form des Bogenelementes einer jeden Fläche zu bestimmen, deren geodätische Curven mehrere infinitesimale Transformationen gestatten.

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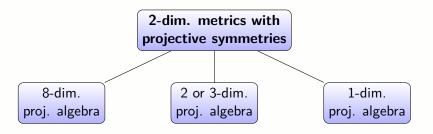
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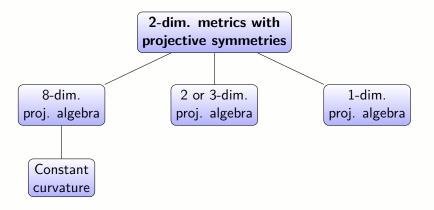
The problem has a natural multi-dimensional generalization, that we shall discuss later

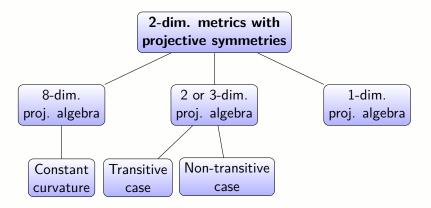
Theorem (Sophus Lie)

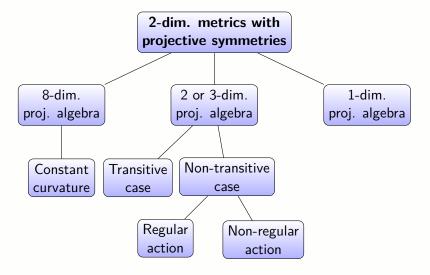
The Lie algebra of projective vector fields on a 2-dimensional manifold can be 0, 1, 2, 3 and 8 dimensional.

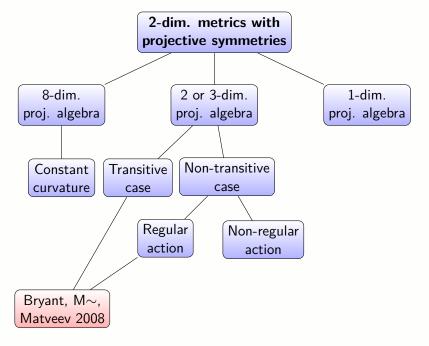
2-dim. metrics with projective symmetries

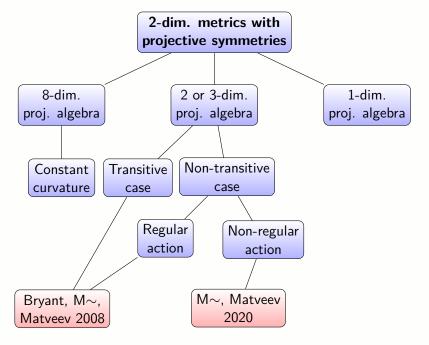


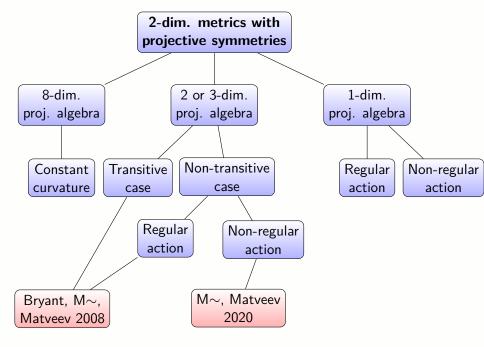


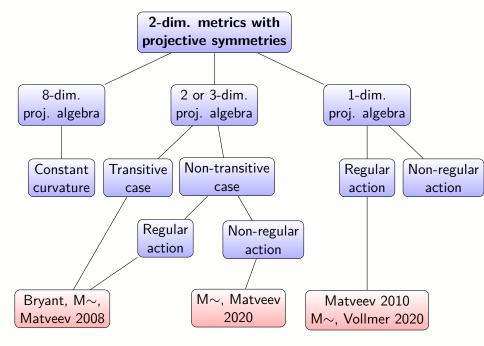


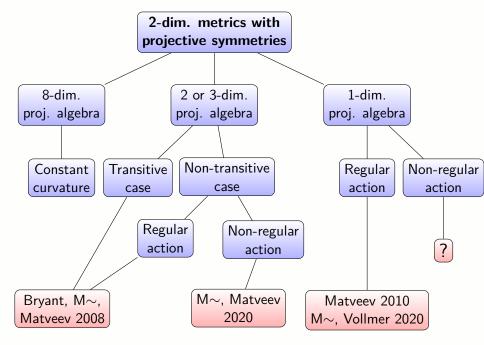


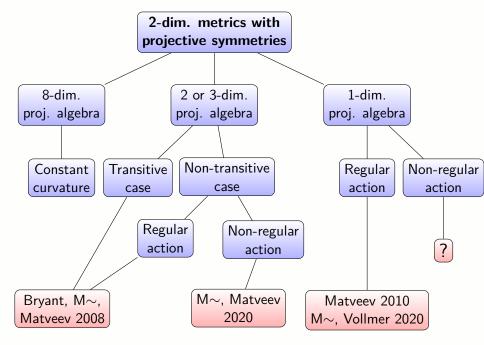












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We shall give more details later.

#### Scheme of the talk

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Equation of geodesics:

$$\ddot{x} + \Gamma_{11}^{1} \dot{x}^{2} + 2\Gamma_{12}^{1} \dot{x} \dot{y} + \Gamma_{22}^{1} \dot{y}^{2} = c(t) \dot{x} \ddot{y} + \Gamma_{21}^{2} \dot{x}^{2} + 2\Gamma_{12}^{2} \dot{x} \dot{y} + \Gamma_{22}^{2} \dot{y}^{2} = c(t) \dot{y}$$

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By eliminating the parameter t we obtain

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But with the substitution  $a = \det(g)^{-\frac{2}{3}}g$ we obtain

$$a_{11x} - \frac{2}{3}F_1 a_{11} + 2F_0 a_{12} = 0$$
  

$$a_{11y} + 2a_{12x} - \frac{4}{3}F_2 a_{11} + \frac{2}{3}F_1 a_{12} + 2F_0 a_{22} = 0$$
  

$$2a_{12y} + a_{22x} - 2F_3 a_{11} - \frac{2}{3}F_2 a_{12} + \frac{4}{3}F_1 a_{22} = 0$$
  

$$a_{22y} - 2F_3 a_{12} + \frac{2}{3}F_2 a_{22} = 0$$

$$y'' = -\Gamma_{11}^2 + (\Gamma_{11}^1 - 2\Gamma_{12}^2)y' - (\Gamma_{22}^2 - 2\Gamma_{12}^1)y'^2 + \Gamma_{22}^1y'^3$$

$$y'' = F_0(x, y) + F_1(x, y)y' + F_2(x, y)y'^2 + F_3(x, y)y'^3$$

$$F_{0} = \frac{1}{2 \det(g)} (g_{12}g_{11x} - 2g_{11}g_{12x} + g_{11}g_{11y})$$

$$F_{1} = \frac{1}{2 \det(g)} (3g_{12}g_{11y} - 2g_{11}g_{22x} + g_{22}g_{11x} - 2g_{12}g_{12x})$$

$$F_{2} = \frac{1}{2 \det(g)} (2g_{12}g_{12y} - 3g_{12}g_{22x} - g_{11}g_{22y} + 2g_{22}g_{11y})$$

$$F_{3} = \frac{1}{2 \det(g)} (2g_{22}g_{12y} - g_{22}g_{22x} - g_{12}g_{22y})$$

$$F_{3} = \det(g)^{-\frac{2}{3}}g$$
we obtain

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 That is linear!

Non-linear!

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Theorem (Eastwood, Matveev (2009))

An N-dim. metric g lies in a given projective class if and only if the following system has solution

$$\nabla_a \sigma^{bc} - \frac{1}{N+1} (\delta^c_a \nabla_i \sigma^{ib} + \delta^b_a \nabla_i \sigma^{ic}) = 0, \ \sigma^{ij} := \det(g)^{\frac{1}{N+1}} g^{ij}$$

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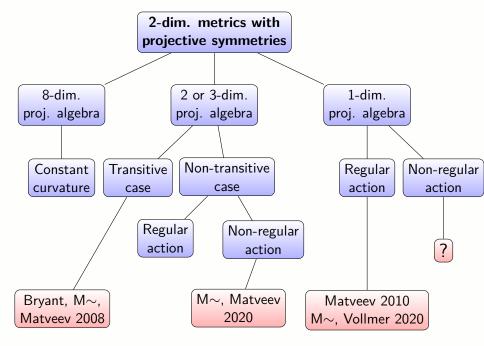
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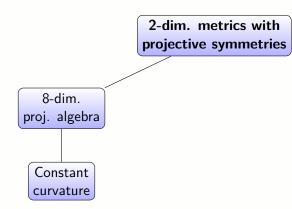
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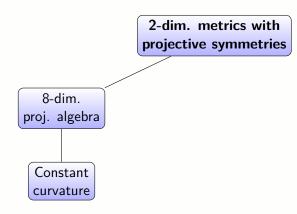
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### Theorem (Bryant, Eastwood, Dunajski (2009))

There exists a differential invariant of order six which decides if a 2-dimensional projective connection is metrizable.

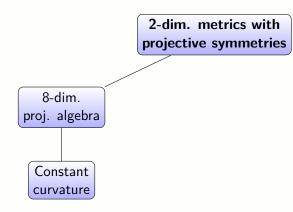






Any second order ODE with 8-dim. Lie symmetry algebra is equivalent to

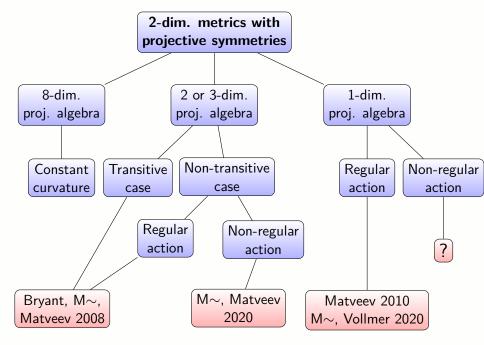
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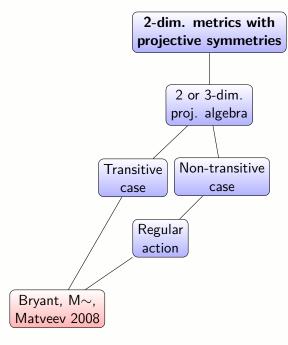


Any second order ODE with 8-dim. Lie symmetry algebra is equivalent to

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and any metric admitting the previous equation as the equation of unparametrized geodesics (projective connection) has constant curvature.





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and by computing the Liouville tensor (Cartan invariants) we realize that they vanish. Thus, (8) is equivalent to y'' = 0, i.e., a metrizable projective connection with metrics of constant curvature.

We now focus on the case when the projective algebra [X, Y] = X is realized as follows:

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#### Theorem (Koenigs)

Let  $\mathcal{I}(g) := \{$ quadratic integrals of the geodesic flow of  $g\}$ . Then dim  $\mathcal{I}(g) \in \{1, 2, 3, 4, 6\}$  and dim  $\mathcal{I}(g) = 6 \longleftrightarrow g$  is of constant curvature.

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#### we obtain the following

#### Theorem

The projective connection (9) comes from a metric iff

$$y'' = C_1 y' \pm e^{-2x} {y'}^3$$

Theorem (R. Bryant, G. Manno, V. Matveev (2008))

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,  $\epsilon_i \in \{-1,1\}$ ,  $b \neq \{-2,0,1\}$   
(b)  $a \left(\epsilon_1 \frac{e^{(b+2)x} dx^2}{(e^{bx}+\epsilon_2)^2} + \frac{e^{bx} dy^2}{e^{bx}+\epsilon_2}\right)$ ,  $a \neq 0$ ,  $\epsilon_i \in \{-1,1\}$ ,  $b \neq \{-2,0,1\}$   
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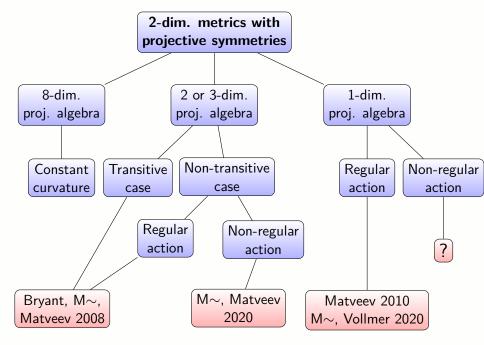
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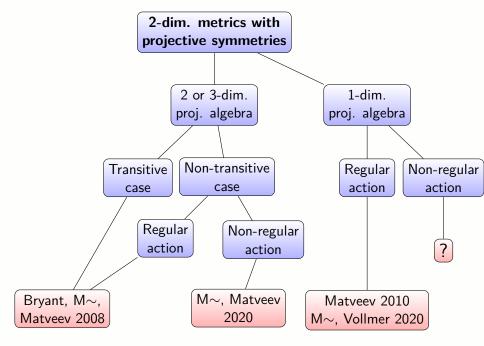
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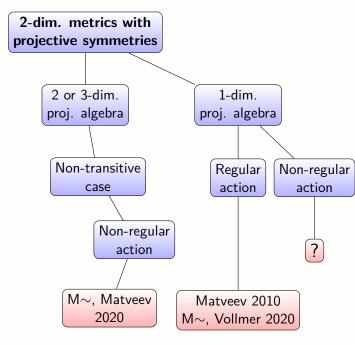
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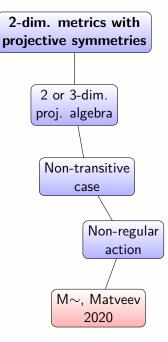
#### 2 projective vector fields $\Rightarrow$ existence of a Killing vector field.

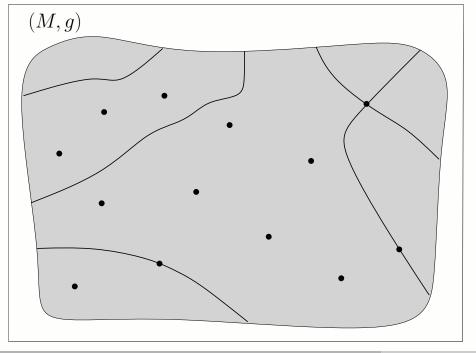
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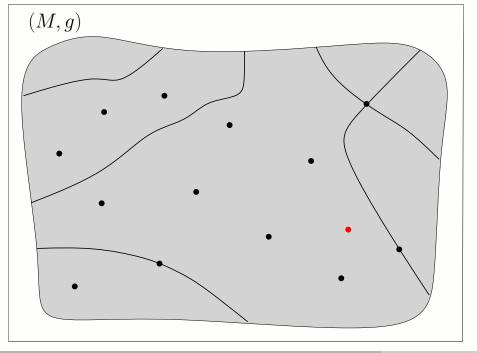




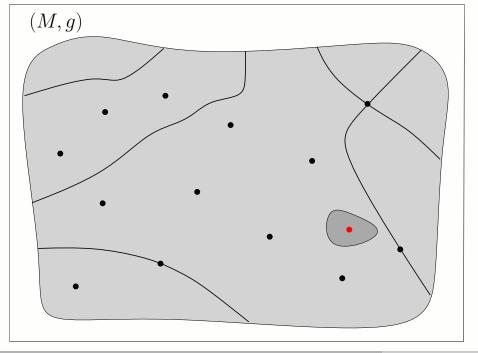


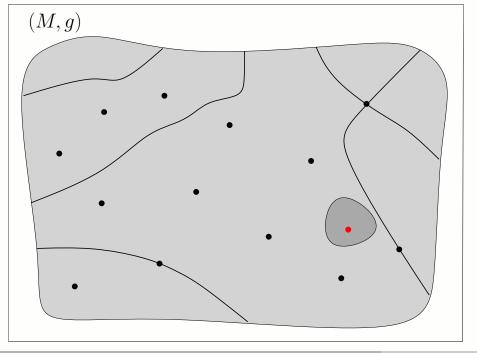


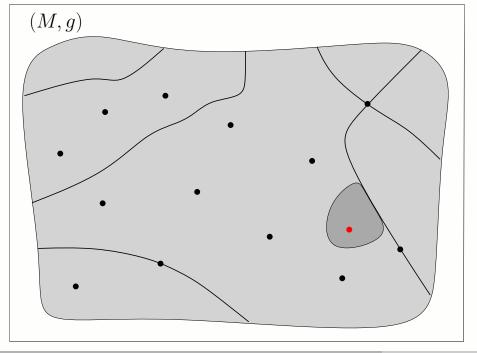


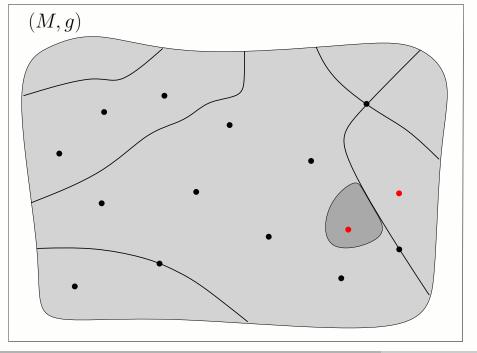


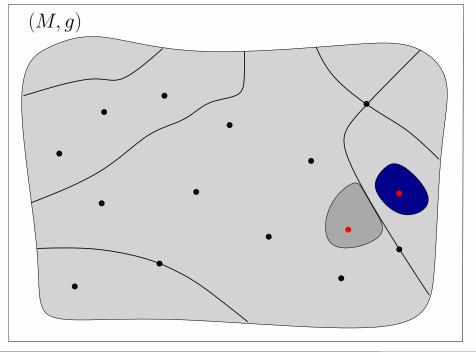
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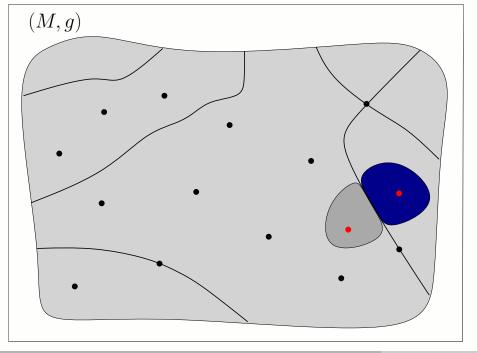




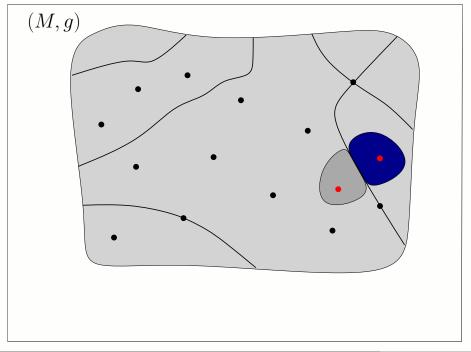


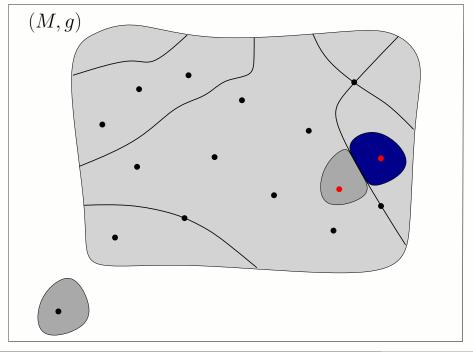


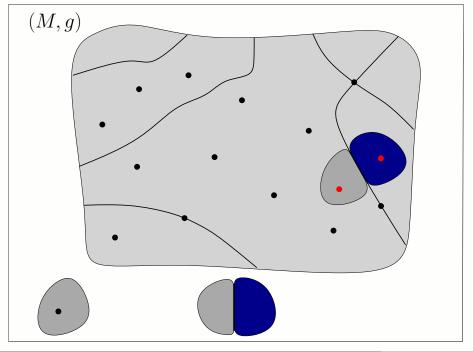


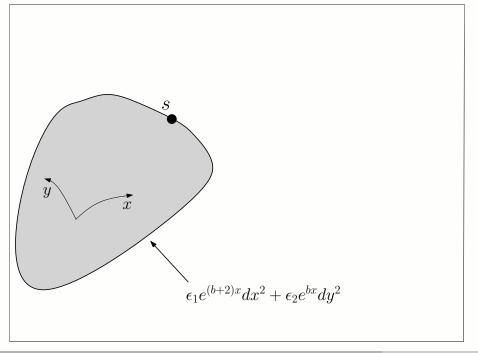


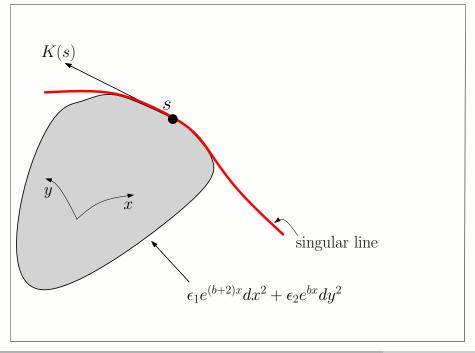
#### Metrics admitting projective symmetries

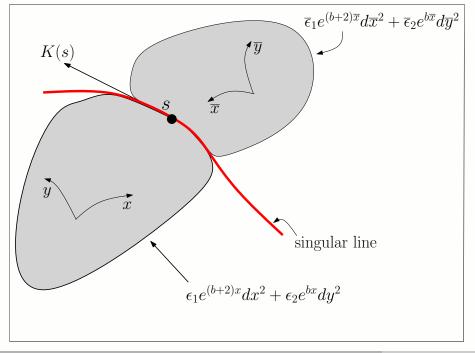


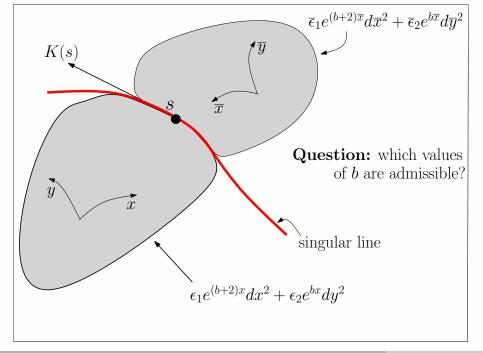


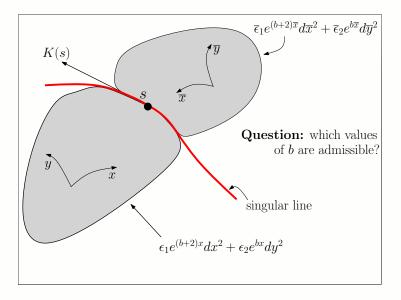


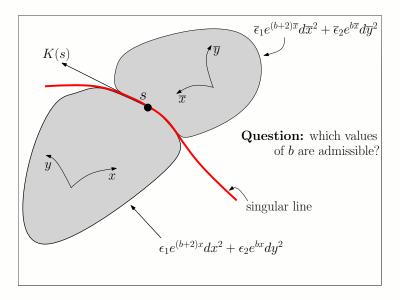




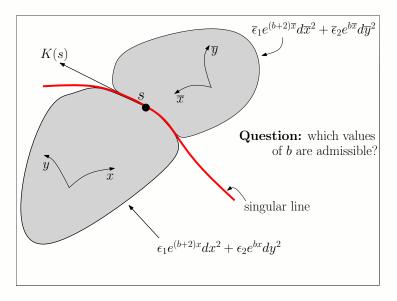




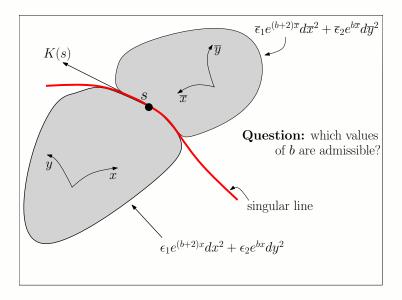


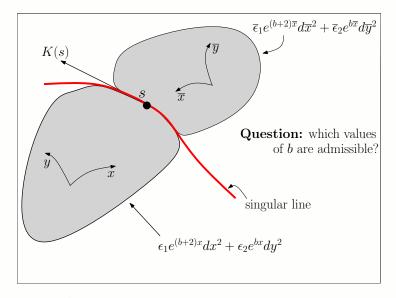


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$$g = (1 + yx^{h})^{-\frac{h+1}{h}} dx dy, \qquad K = x^{h+1} \frac{\partial}{\partial x} + h \frac{\partial}{\partial y}, \quad H = -x \frac{\partial}{\partial x} + hy \frac{\partial}{\partial y}$$

**1** Metrics admitting precisely two projective vector fields

$$\begin{array}{ll} (\mathsf{A}) & (1+yx^h)^{-\frac{h+1}{h}}dxdy, & h \in \mathbb{N} \setminus \{1\}; \\ (\mathsf{B}) & \frac{1}{(1+\epsilon_2x^h)^2}dx^2 + \epsilon_1 \frac{1}{1+\epsilon_2x^h}dy^2, & \epsilon_i \in \{-1,1\}, & h \in \mathbb{N} \setminus \{1,2\}. \end{array}$$

- 2 Metrics admitting precisely three projective vector fields

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**2** Metrics admitting precisely three projective vector fields

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$$(1 + x^2 + \epsilon y^2)(dx^2 + \epsilon dy^2), \qquad \epsilon \in \{-1, 1\}$$
  
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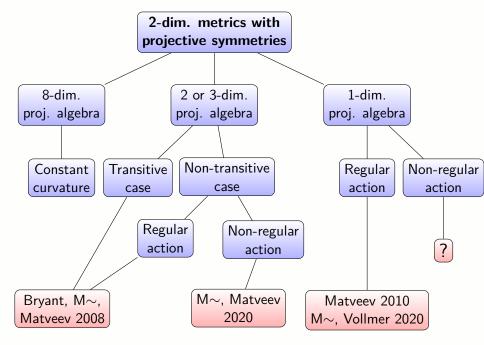
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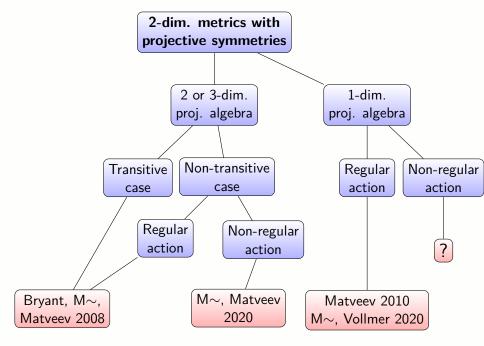
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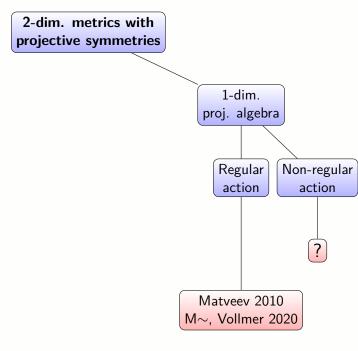
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2 projective vector fields  $\Rightarrow$  existence of a Killing vector field.







We have already seen that for solving such problem it is crucial to solve the "Metrizability Equations"

$$\nabla_a \sigma^{bc} - \frac{1}{3} (\delta^c_a \nabla_i \sigma^{ib} + \delta^b_a \nabla_i \sigma^{ic}) = 0, \quad \sigma^{ij} := \det(g)^{\frac{1}{3}} g^{ij}. \tag{10}$$

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**Key observation:** The flow of (the unique up to a constant) projective vector field send a metric  $g[K_1, K_2]$  in one of the same form. We can then distinguish non-isometric metrics by studying some invariants, for instance the length of the projective vector field along its flow.

#### Example

The metric

$$g = k \frac{y - x}{xy} \left( \frac{e^{-3x}}{x} dx^2 + h \frac{e^{-3y}}{y} dy^2 \right)$$

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The list is much more long, see

G. MANNO, A. VOLLMER: Normal forms of two-dimensional metrics admitting exactly one essential projective vector field, *J. Math. Pure Appl.*, **135** (2020), 26–82.

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$$g[K_1, K_2] = \frac{\left(K_1 \det(g_1)^{\frac{1}{4}} g_1^{-1} + K_2 \det(g_2)^{\frac{1}{4}} g_2^{-1}\right)^{-1}}{\det\left(K_1 \det(g_1)^{\frac{1}{4}} g_1^{-1} + K_2 \det(g_2)^{\frac{1}{4}} g_2^{-1}\right)}$$

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For metrics  $g_1$  and  $g_2$  such that the (1,1)-tensor

$$\left.\frac{\det(g_2)}{\det(g_1)}\right|^{\frac{1}{4}}g_2^{-1}g_1$$

is diagonalizable, we have a simple local description.

Metrics admitting projective symmetries

Metrics  $g_1$  and  $g_2$  assume either the form

$$g_{1} = \pm (F_{1} - F_{2})(F_{1} - F_{3}) (dx^{1})^{2} \pm (F_{2} - F_{1})(F_{2} - F_{3}) (dx^{2})^{2} \pm (F_{3} - F_{1})(F_{3} - F_{2}) (dx^{3})^{2} g_{2} = \pm \frac{(F_{1} - F_{2})(F_{1} - F_{3})}{F_{1}^{2}F_{2}F_{3}} (dx^{1})^{2} \pm \frac{(F_{2} - F_{1})(F_{2} - F_{3})}{F_{1}F_{2}^{2}F_{3}} (dx^{2})^{2} \pm \frac{(F_{3} - F_{1})(F_{3} - F_{2})}{F_{1}F_{2}F_{3}^{2}} (dx^{3})^{2}$$

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, or the form  
 $g_1 = \zeta(z) (h \pm dz^2)$   $g_2 = \frac{\zeta(z)}{Z(z) \rho^2} \left(\frac{h}{\rho} \pm \frac{dz^2}{Z(z)}\right)$ 

where  $h = h_{11}dx^2 + 2h_{12}dxdy + h_{22}dy^2$ ,  $h_{ij} = h_{ij}(x, y)$ , and

$$\zeta(z) = Z(z) - 
ho, \quad 
ho \in \mathbb{R}.$$

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The metric, for  $k_i \neq 0$ ,

$$g = k_1 \left(\frac{1}{x} - \frac{1}{y}\right) \left(\frac{1}{x} - \frac{1}{z}\right) e^{2x} dx^2 + k_2 \left(\frac{1}{y} - \frac{1}{x}\right) \left(\frac{1}{y} - \frac{1}{z}\right) e^{2y} dy^2$$
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The list of course is much longer, see G. MANNO, A. VOLLMER: 3-dimensional Levi-Civita metrics with projective vector fields, https://arxiv.org/abs/2110.06785

## Complex case

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More precisely, if  $(M, J, \nabla)$  is a real 2*n*-dimensional smooth manifold equipped with a complex structure J and a complex connection  $\nabla$ , i.e., a torsion free affine connection such that  $\nabla J = 0$ , a *J*-planar curve is a regular curve  $\gamma : I \subseteq \mathbb{R} \to M$  such that

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Promising results have been obtained with Jan Schumm and Andreas Vollmer.