Conformal geometric aspects of hyperplane sections of Lagrangian Grassmannians

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ongoing work with Jan Gutt and Giovanni Moreno October 23, 2015

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$$f(x, y, u, u_x, u_y) = 0$$
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From now we fix the point  $m \in M$ , i.e. we work in a fibre.















Hyperplane sections are Monge-Ampère equations:

$$a_0 + a_1 u_{xx} + a_2 u_{xy} + a_3 u_{yy} + a_4 (u_{xx} u_{yy} - u_{xy}^2) = 0$$
  
Plücker :  $(p_{11}, p_{12}, p_{22}) \hookrightarrow (1, p_{11}, p_{12}, p_{22}, p_{11} p_{22} - p_{12}^2)$ 

Now we see how to define Monge-Ampère equations by using characteristics directions and how to characterize such PDEs in terms of some canonical conformal structures on Lagrangian Grassmannians



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Locally

$$\mathcal{E}_{\mathcal{D}} : \det \left( \begin{array}{cc} u_{xx} - f_{11} & u_{xy} - f_{12} \\ u_{xy} - f_{21} & u_{yy} - f_{22} \end{array} \right) = 0$$

 $\mathcal{E}_{\mathcal{D}}$  is a parabolic una Monge-Ampère equation  $\iff f_{12} = f_{21} \iff \mathcal{D} = \mathcal{D}^{\perp}$ .

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So, we have a natural conformal structure on LGr(2, 4)

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• A straightforward computation shows that

 $II - H \cdot I$ , H mean curvature

is a conformal invariant.

 $S \subset LGr(2,4)$  is a hyperplane section  $\iff$  II – H · I = 0

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Let us consider the symmetric 3-tensor T:

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Hyperquadric sections?